

MoEDAL, MAPP and future endeavours

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The unprecedented collision energy of the LHC has opened up a new discovery frontier, however, without any signs of new physics in sight. The first LHC dedicated search experiment, MoEDAL, started data taking for Run 2. MoEDAL is designed to search highly ionising particle avatars of new physics using pp and heavy-ion collisions at the LHC. The planned upgrade for MoEDAL at Run 3 of the LHC — the MAPP detector (MoEDAL Apparatus for Penetrating Particles) — will extend MoEDAL's physics reach to include feebly interacting, long-lived messengers of physics beyond the Standard Model. This will allow the experiment to explore a number of models of new physics, including dark sector models, in a complementary way to that of conventional LHC collider experiment detectors. Furthermore, a possible astroparticle extension to MoEDAL, called Cosmic-MoEDAL, will allow the search for magnetic monopoles to be continued from the TeV scale to the GUT scale. This paper focuses on recent results and plans for the LHC Run 3.

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1. Introduction

The MoEDAL (Monopole and Exotics Detector at the LHC) [1] experiment at the Large Hadron Collider (LHC) [2] is dedicated to searches for manifestations of new physics through highly ionising particles in a manner complementary to ATLAS and CMS [3]. It is the first dedicated *search* LHC experiment to be approved among others that followed [4–6]. The principal motivation for MoEDAL is the quest for magnetic monopoles, as well as for any massive, stable or long-lived (LL), slow-moving particle with the fundamental electric charge (or multiples thereof) arising in various extensions of the Standard Model (SM) [7], such as supersymmetric partners [8] and D-matter [9–15], among others [7, 16]. Emphasis is given here on recent MoEDAL results, based on the exposure of magnetic monopole trapping volumes to pp and heavy-ion collisions, on future prospects, including electric charges, and on the description and sensitivity of the planned MoEDAL Apparatus for Penetrating Particles (MAPP), designed to extend the LHC reach in the quest for dark matter [17, 18].

2. The MoEDAL experiment

The MoEDAL detector [19] is deployed around the intersection region at Point 8 (IP8) of the LHC in the LHCb Vertex Locator cavern. It is a unique and largely passive detector based on three different detection techniques.

A large array of CR39®, Makrofol® and Lexan® nuclear track detector (NTD) stacks surround the IP8. The passage of a highly ionising particle through the plastic sheets is marked by an invisible damage zone along the trajectory. The damage zone is revealed as a cone-shaped etch-pit when the plastic detector is chemically etched in the INFN Bologna laboratory. Then the sheets of plastics are scanned looking for aligned etch pits in multiple sheets.

A unique feature is the use of magnetic monopole trappers (MMTs) to capture magnetically charged particles. The aluminium absorbers of MMTs are subject to an analysis looking for magnetically charged particles at the ETH Zurich SQUID laboratory [20].

The only active MoEDAL sub-detector comprises an array of several TimePix pixel devices dedicated to the monitoring of cavern background sources. The time-over-threshold mode allows a 3D mapping of the charge spreading effect in the silicon sensor volume [21].

3. Searches for monopoles and dyons in MoEDAL

MoEDAL is designed to fully exploit the energy-loss mechanisms of magnetically charged particles [22–25] in order to optimise its potential to discover these elusive particles. Various theoretical scenarios foresee the production of magnetic charge at the LHC [7, 26]: (light) 't Hooft-Polyakov monopoles [24, 25, 27], electroweak monopoles [28–33], global monopoles [34–40], monopoles in Born-Infeld theory [41–44] and monopolium [23, 45–49], a monopole-pair bound state. Magnetic monopoles and dyons — the latter possessing both magnetic and electric charge [50] — are fascinating hypothetical particles. Even though there is no empirical evidence for their presence, strong theoretical reasons motivate their existence and many theories, including grand unified theories [24, 25] and superstring theory [51, 52], predict their existence.

Until recently, the MoEDAL results on monopoles were based on the scanning of the MMTs, exposed to LHC Run 1 8-TeV data [53] and to 13 TeV pp collisions [54–56]. The SQUID analysis yielded no observed isolated magnetic charges, leading to upper limits on monopole production cross sections. This outcome led to lower mass exclusion bounds when considering two pair production processes: (a) a Drell-Yan-like (DY) process in photon s -channel intermediation, and (b) a photon-fusion t -channel diagram [57].

If the DY production mechanism is considered, the ATLAS bounds [58–60] are better than the MoEDAL ones for $|g| \leq 2g_D$ due to the higher luminosity delivered in ATLAS and the loss of acceptance in MoEDAL for small magnetic charges, while MoEDAL is the only detector sensitive to high charges. The production cross section at the LHC energies for photon fusion is much higher than the DY [57], so MoEDAL, being the only experiment considering it, set the most stringent limits on monopoles overall [61].

MoEDAL performed recently the first dedicated dyon search in a collider experiment by means of MMT scanning; a summary of the dyon mass limits are shown in Figure 1. Mass limits in the range 750–1910 GeV were set using a benchmark DY production model for dyons with magnetic charge up to $5g_D$, for electric charge from $1e$ to $200e$, and for spins 0, $1/2$ and 1 [62]. Moreover, an analysis for trapped monopoles in the Run 1 CMS beam pipe is currently underway [63].

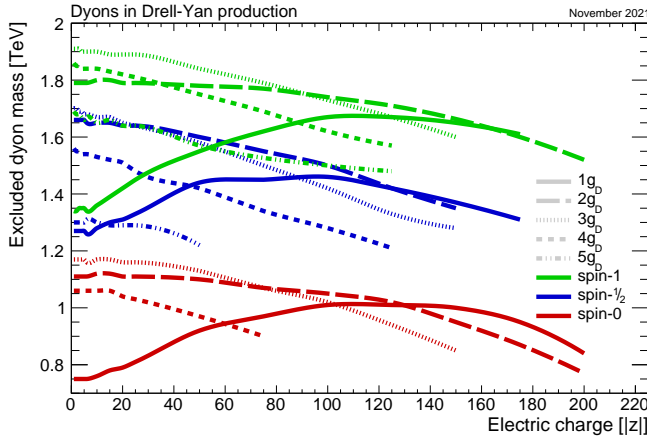


Figure 1: Dyon mass limits obtained by MoEDAL [62] at $\sqrt{s} = 13$ TeV as a function of electric charge for various spins and magnetic charges, assuming Drell-Yan pair-production mechanism.

In both production processes, the monopole pair couples to the photon via a coupling that depends on g_D and therefore has a value of $O(10)$. This large monopole–photon coupling invalidates any perturbative treatment of the cross-section calculation and hence any result based on it is *only indicative* and used merely to facilitate comparisons between experiments. However, it is stressed that the upper bounds placed on production cross sections are solid and can be relied upon.

This situation is resolved if thermal Schwinger production of monopoles in heavy-ion collisions is considered [64]. This mechanism becomes effective in the presence of strong magnetic fields and calculations rely on semiclassical techniques rather than perturbation theory, thus overcoming these limitations [65–70]. Heavy-ion collisions at the LHC produce the strongest known magnetic fields in the current Universe, and the first search for such production was conducted by MoEDAL during the 5.02 TeV/nucleon heavy-ion run, during which the MMTs were exposed to 0.235 nb^{-1} of Pb-Pb collisions and analysed later with a SQUID. Monopoles with Dirac charges $1g_D \leq g \leq 3g_D$

and masses up to 75 GeV were excluded, as seen in Figure 2. This analysis provided the first lower mass limits for finite-size monopoles from a collider search [71].

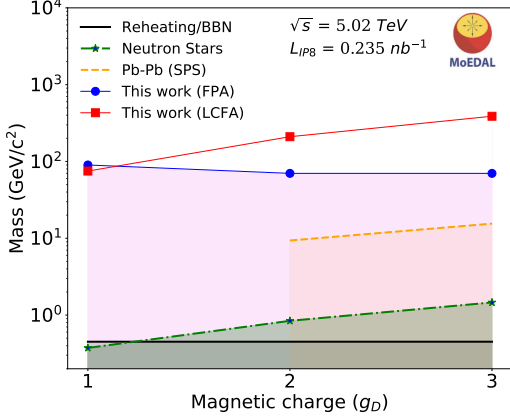


Figure 2: Magnetic-monopole 95% CL exclusion region in mass and charge obtained using two different calculations of Schwinger production mechanism for Pb-Pb collisions. From [71].

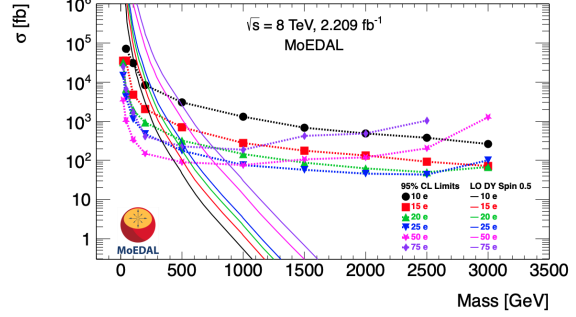


Figure 3: Cross-section upper limits (at 95% CL) for monopoles produced with a Drell-Yan pair production mechanism in 13 TeV pp collisions as a function of mass for spin- $1/2$ HECOs with electric charge $10e-75e$. The solid lines are cross-section calculations at leading order. From [72].

4. Electrically charged particles

The MoEDAL detector is also designed to search for any massive, long-lived, slow-moving particles with single or multiple electric charge arising in many scenarios of physics beyond the Standard Model. Supersymmetric (SUSY) LL particles [16, 73–76], quirks, strangelets, Q-balls, and many others fall into this category [7]. A generic search for high-electric-charge objects (HECOs) using for the first time NTD data has been completed recently [72] with a prototype NTD detector exposed at 8-TeV pp collisions. Upper cross section limits on fermionic HECOs are shown in Figure 3. The limits placed on the DY production cross sections of HECO pairs vary from ~ 30 fb to 70 pb, for electric charges in the range $15e$ to $175e$ and masses from 110 GeV to 1020 GeV. For comparison, ATLAS has constrained HECOs of electric charge between $20e$ to $100e$ [60].

MoEDAL is mostly sensitive to slow-moving LL superpartners unlike ATLAS/CMS suitability for faster ones ($\beta \gtrsim 0.8$) [77]. However, the lower integrated luminosity for MoEDAL at IP8 remains a limiting factor for simple scenarios. Direct production of heavy fermions abundantly produced via the strong interaction is the most favourable scenario for MoEDAL. Complex topologies appear to be promising for LL sleptons in phenomenologically realistic models, where MoEDAL could cover parameter space less accessible by CMS [75] and ATLAS [76]. Even for SUSY models observable by both ATLAS/CMS and MoEDAL, the latter’s added value remains, thanks to the completely different detector and analysis techniques, involving uncorrelated systematic uncertainties.

The prospects for detecting particles of higher electric charges in MoEDAL are also very promising. Doubly charged scalars and fermions are suggested by Type-II ($H^{\pm\pm}$) and Type-III seesaw models of neutrino masses, respectively [78]. In addition to the Type-II seesaw model, several other scenarios, namely the Left-Right model, the Georgi-Machacek model, the 3-3-1 model and the little-Higgs model also predict doubly charged scalars. Even better discovery reach

is anticipated for charges of $2e$, $3e$ and $4e$, proposed in radiative neutrino mass models, which often add a discrete symmetry to the SM gauge group [78]. For such models, at least one signal event at the NTDs is expected for up to masses of 290, 610 and 960 GeV for scalars $S^{\pm 2}$, $S^{\pm 3}$ and $S^{\pm 4}$ in Run 3 [79]. Recent studies [80, 81] quantified the MoEDAL potential to discover generic electrically charged scalars and fermions in the range $1e$ to $6e$ in the High Luminosity LHC (HL-LHC) runs with sensitivity superior to that of ATLAS and CMS.

5. MoEDAL Apparatus for Penetrating Particles

MoEDAL proposes to deploy MAPP in a gallery near IP8 shielded by an overburden of approximately 100 m of limestone from cosmic rays [82, 83]. The Phase-1 MAPP sub-detector [84] is composed of mQP, a detector sensitive to particles with fractional charge as small as $0.001e$, the so-called millicharged particles, mCP. MAPP-mQP is made of plastic scintillation bars, as shown in Figure 4, and it is currently been installed in the UA82 gallery at a distance 100 m from IP8. It is expected to start taking data in 2023.

Another part of the detector, the MAPP-LLP, will be deployed in the Phase-2 MAPP as three nested boxes of scintillator hodoscope detectors, in a ‘Russian doll’ configuration, following as far as possible the contours of the UG1 cavern. It is designed to be sensitive to LL neutral particles from new physics scenarios via their interaction or decay in flight in a decay zone of size approximately 5 m (wide) \times 10 m (deep) \times 3 m (high). The MAPP detector can be deployed in a number of positions in the forward direction, at a distance of $O(50\text{ m})$ from IP8. It will be installed during the Long Shutdown 3 to be operated in the HL-LHC [85].

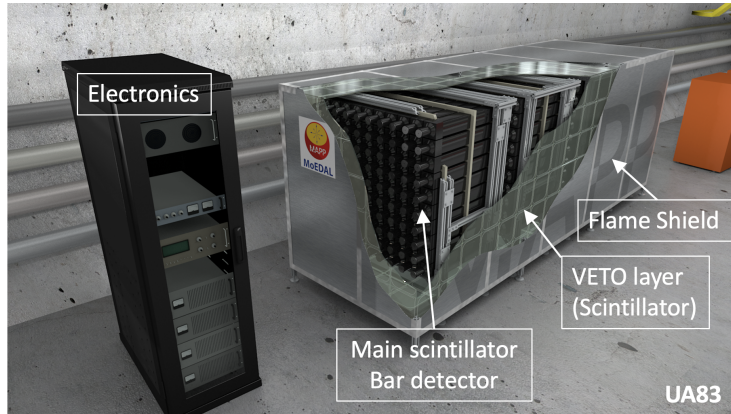


Figure 4: A schematic view of the MAPP-mQP subdetector.

The mCPs arise when a new $U(1)$ is introduced to a dark sector with a massless dark photon A' , coupled to the SM photon field, and a massive dark fermion ψ with charge much less than that of an electron as a result of kinetic mixing [86]. MAPP-mQP is expected to significantly extend the reach obtained by milliQan demonstrator [87]. MAPP-mQP can detect a heavy neutrino with large electric dipole moment considered to be a member of a fourth generation lepton doublet [88].

The MoEDAL-LLP detector should be sensitive to portal interactions that connect a hidden (dark) sector and the visible sector of the SM. Scenarios beyond-the-SM that introduce a dark sector in addition to the visible SM sector are required to explain a number of observed phenomena in

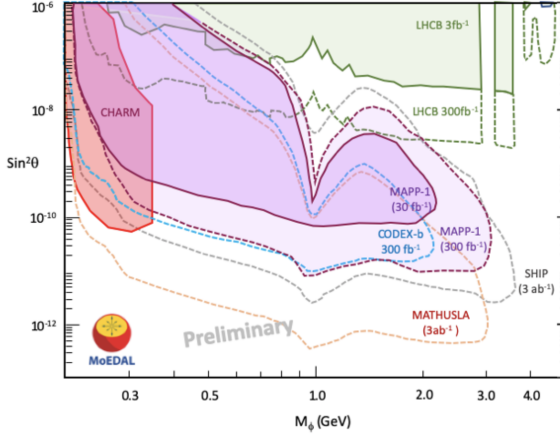


Figure 5: MAPP-LLP sensitivity with 30 and 300 fb⁻¹ of pp collisions for a dark Higgs boson ϕ in terms of its mass and its coupling to SM particles.

particle physics, astrophysics and cosmology such as the non-zero neutrino masses and oscillations, the dark matter, baryon asymmetry of the Universe, the cosmological inflation. In particular, dark Higgs bosons interact with the SM through a kinetic mixing term, thus probing one of the few possible renormalisable interactions with a hidden sector, the Higgs portal quartic scalar interaction. Such scenarios are accessible to MAPP and other future experiments [4–6]. A comparison of the projected sensitivity of MAPP-LLP with other LLP experiments is illustrated in Figure 5. Regarding their cosmological implications, dark Higgs bosons may mediate interactions with hidden dark matter that has the correct thermal relic density or resolves small-scale-structure discrepancies [89].

In the fermion portal, right-handed LL heavy neutrinos can be pair produced in the decay of an additional Z' boson in the gauged $B - L$ model, which also contains a singlet scalar field that spontaneously breaks the extra $U(1)_{B-L}$ gauge symmetry [90]. In this model, MAPP will fill the gap left by other LHC experiments [91]. Sterile neutrinos may be LL in neutrino-extended SM effective field theories, ν SMEFT, and can be produced in leptonic and semi-leptonic decays of charmed and bottomed mesons, decaying to leptons via neutral and charged weak currents, thus becoming detectable in MAPP-LLP [92]. R -parity violating supersymmetry also predicts LLPs, such as light LL neutralinos $\tilde{\chi}_1^0$ decaying via λ'_{ijk} couplings to charged particles. In benchmark scenarios with charm or bottom mesons decaying into $\tilde{\chi}_1^0$, MAPP can cover various $\tilde{\chi}_1^0$ lifetimes [93], in similar fashion as in sterile neutrinos.

6. The Cosmic-MoEDAL

The MoEDAL collaboration is considering [82, 94] an astroparticle extension to the MoEDAL experiment that will enable the search for magnetic monopoles of mass of up to the Grand Unification (GUT) scale. The detector technology for the “Cosmic-MoEDAL” will be the same as its LHC counterpart, namely based on NTDs. SLIM, the first experiment to use this approach to search for highly ionising particles of cosmic origin [95, 96], was deployed at high altitude at the Mt Chacaltaya laboratory in Bolivia with an elevation of 5,400 m. However, its modest size (400 m²) precluded it from the search for a flux of cosmic monopoles below the Parker Bound.

Cosmic-MoEDAL is proposed as a 10,000 m² array of plastic NTDs (CR39®) deployed at high altitude. Such an array would be able to take the search for cosmic monopoles with velocities

$\beta \gtrsim 0.1$ from the TeV scale to the GUT scale for monopole fluxes well below the Parker Bound, as shown in Figure 6.

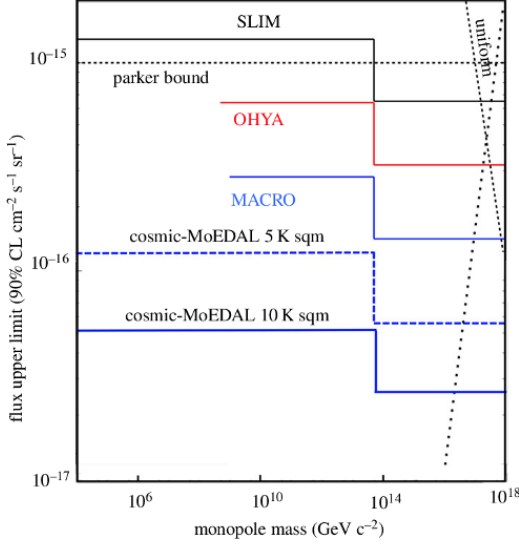


Figure 6: Flux upper limits for cosmic magnetic monopoles of charge $1g_D$ and $\beta > 0.05$ versus monopole mass. The figure shows the 90% CL limits obtained by the SLIM [95], MACRO [97] and OHYA [98] experiments. From [82].

7. Summary and outlook

MoEDAL, the first dedicated *search* experiment at the LHC, is extending considerably the LHC reach in the search for (meta)stable highly ionising particles, predicted in a variety of theoretical models including magnetic monopoles, SUSY long-lived superpartners, D-matter, quirks, strangelets, Q-balls, among others. MoEDAL is optimised to probe precisely all such LL states, unlike the other LHC experiments, by combining different detector technologies: plastic nuclear track detectors, trapping volumes and pixel sensors. It provides the best sensitivity to high magnetic charges. The latest highlights in MoEDAL results include the first search for dyons at the LHC, the first search for monopoles produced via the Schwinger mechanism and the first search with NTDs constraining electrically charged LL particles.

The search for trapped monopoles in the Run 1 CMS beam pipe is expected soon. MoEDAL will be sensitive to single and multiple electric charges predicted in supersymmetric models and other exotic scenarios in LHC Run 3 [99]. Moreover, the search for HIPs such as Q-balls, nuclearites and magnetic monopoles can be pushed to GUT mass scales, and below the Parker Bound, by the deployment of the Cosmic-MoEDAL array at high altitude. The operation of the MAPP sub-detector will allow MoEDAL's physics reach to be significantly expanded to include the quest of millicharged and weakly interacting LL neutral particles, arising in hidden sectors.

Acknowledgments

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