

## Highlights from the ALICE collaboration

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#### Abstract

Significant advances in heavy-ion physics have come about as almost ten years have passed since the beginning of the heavy-ion programme at the LHC. In this document, we review some key results from the ALICE collaboration on a variety of topics, from properties of the quark-gluon plasma (QGP) to basic QCD, and discuss how these set the stage for further measurements in the next decade.

*40th International Conference on High Energy physics - ICHEP2020*

*July 28 - August 6, 2020*

*Prague, Czech Republic (virtual meeting)*

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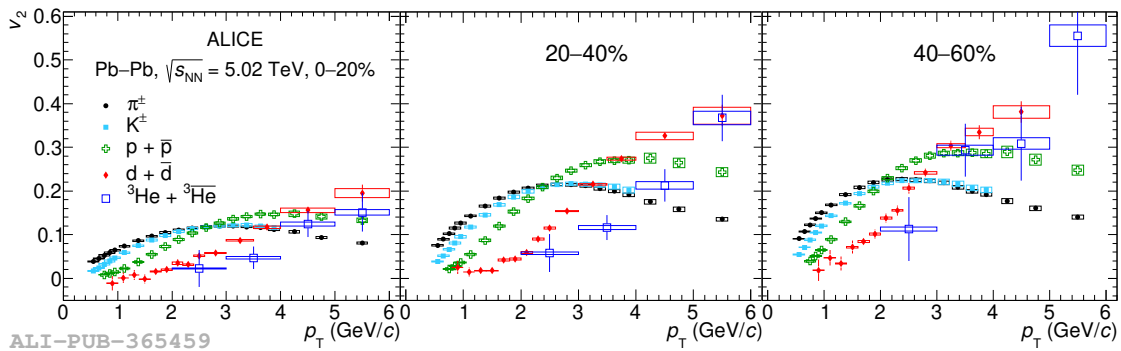
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## 1. Introduction

The primary mission of the ALICE collaboration is the study of the properties of the quark-gluon plasma (QGP) created in relativistic heavy-ion collisions [1]. In doing so, the collaboration accesses matter whose behaviour is dictated by QCD at extremely high densities, similar to those thought to have been at play during the very early stages of the universe. Several phenomena appear in such conditions and have been systematically probed by the collaboration, leading to remarkable advances in our understanding of the QGP. These advances can be broadly classified into three broad categories: (a) the study of the basic properties of the QGP, (b) the determination of under which conditions high-density signatures such as collective behaviour and strangeness enhancement develop, which includes also measurements performed in proton-proton (pp) and proton-lead (p-Pb) collisions and (c) employing the outstanding particle identification capabilities of ALICE to explore many other QCD topics and phenomena. These three classes of studies will be discussed in the next sections.

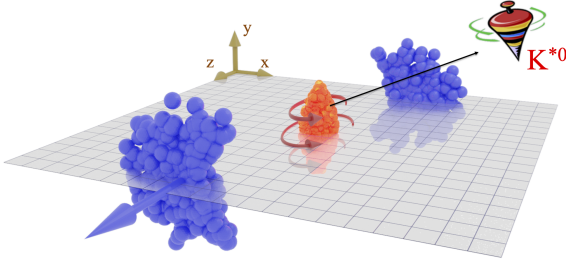
## 2. Properties of the quark-gluon plasma

A fundamental property of the quark-gluon plasma is that it conveys a strong collective motion broadly denoted as ‘flow’ to the emitted hadrons. This phenomenon has been studied in great detail using Fourier decompositions of the azimuthal particle distribution. The leading non-trivial component, the elliptic flow  $v_2$ , has been observed in Pb-Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV for a variety of light-flavour particles, ranging from  $\pi$  to heavier objects such as deuterons or  ${}^3\text{He}$  [2–4], as can be seen in Fig. 1. In all cases,  $v_2$  has been shown to exhibit mass ordering at low momenta, which is expected from hydrodynamic models [5]. Moreover, the  $v_2$  of hadrons containing heavy-flavour quarks has also been measured and has been observed to be non-zero for D mesons,  $J/\psi$  and even beauty decay electrons [6, 7]. Heavier states such as the  $Y(1S)$ , however, exhibit  $v_2$  consistent with zero within uncertainties in the measured momentum range [8], which indicates that either they do not participate in the collective motion or they do so only outside of the measured range due to their very large masses.

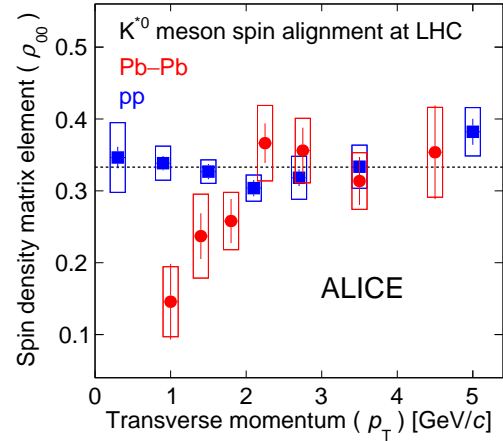


**Figure 1:** Elliptic flow as a function of  $p_T$  for  $\pi$ , K, p, d and  ${}^3\text{He}$  [2, 3].

The motion of the system not only confers linear momentum but also angular momentum to the emitted hadrons. This is a consequence of the fact that in a semi-central collision, the system



**Figure 2:** Illustration of a rotating quark-gluon plasma emitting a  $K_0^*$  meson with its spin aligned with the rotation of the emitting medium.

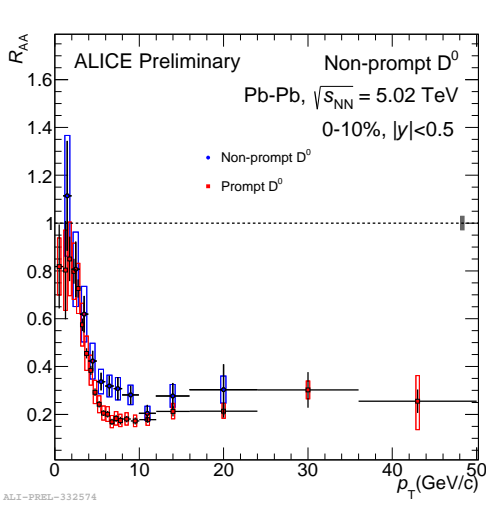


**Figure 3:**  $K_0^*$  spin density matrix element  $\rho_{00}$  as a function of transverse momentum measured in Pb-Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV and in proton-proton collisions at  $\sqrt{s} = 13$  TeV. Adapted from [9].

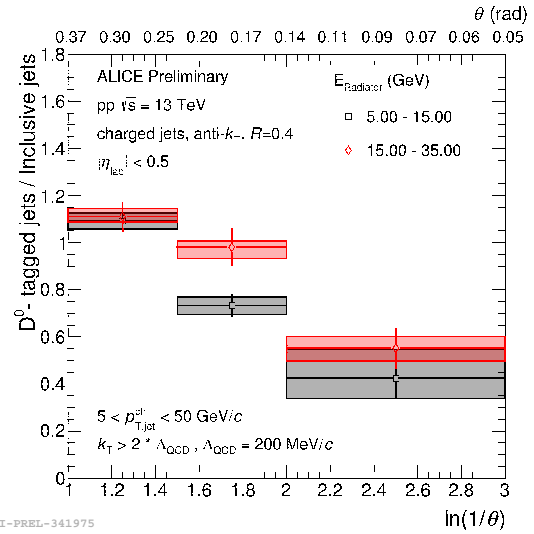
as a whole rotates, as indicated in Fig. 2. As a consequence, the spin of vector mesons such as  $K^{*0}$  is expected to be aligned with that motion, which is what is observed in measurements of the spin density matrix element  $\rho_{00}$  in Pb-Pb collisions in Fig. 3 [9, 10]. It is worth noting that the magnitude of the effect is significantly more pronounced than the  $\Lambda$  polarization measurements by ALICE and STAR [11, 12], which poses an as-of-yet unresolved challenge to theorists.

The strongly interacting medium also blocks the passage of very-high momentum particles, leading to jet quenching. This effect can be quantified via the nuclear modification factor  $R_{AA}$ , which is the ratio between the transverse momentum distributions observed in Pb-Pb and an appropriately scaled spectrum measured in pp collisions at the same centre-of-mass energy. The  $R_{AA}$  has now been notably measured also for prompt  $D^0$  as well as non-prompt  $D^0$  mesons, with the latter coming from the decay of beauty hadrons [13]. While in both cases a strong suppression is observed in the  $R_{AA}$ , as seen in Fig. 4, the non-prompt  $D^0$  nuclear modification factor is slightly higher than the prompt  $D^0$  value between a transverse momentum of 5 and 10 GeV/c [8], suggesting that beauty quarks undergo less quenching in the QGP than charm quarks. A potential explanation for this is the dead cone effect, according to which collinear gluon emissions are suppressed in  $\theta < m_q/E_q$ , where  $\theta$  is the emission angle and  $m_q$  and  $E_q$  are the mass and energy of the emitting quark, respectively.

A more direct consequence of the dead cone effect has, in fact, also been observed in pp collisions at  $\sqrt{s} = 13$  TeV, i.e. for the case of vacuum radiation. More specifically, the observable consequence of this effect is that charged-particle yield at small angles with respect to the jet axis are expected to be suppressed if the jet originated from a heavy-flavour quark. This effect can be seen very clearly in Fig. 5, in which the ratio of jet-associated particle yields is plotted as a function of the logarithm of the inverse of the opening angle for inclusive and for  $D^0$ -tagged jets. It is visible that, at small opening angle  $\theta$  (large values of  $\ln(1/\theta)$ ),  $D^0$ -tagged jets only exhibit half of the charged-particle yields than an inclusive jet sample [14].



**Figure 4:** Nuclear modification factor  $R_{AA}$  as a function of transverse momentum for prompt and non-prompt  $D^0$  mesons in Pb-Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV.



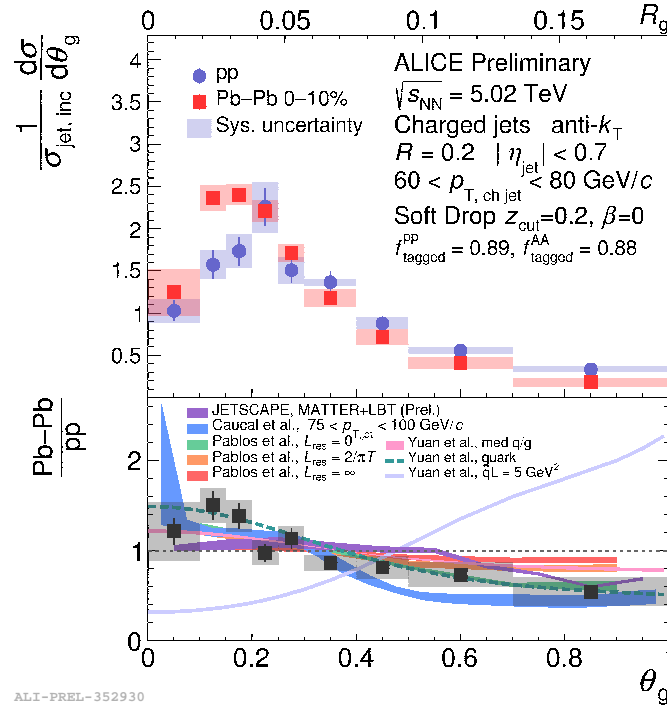
**Figure 5:** Ratio of charged-particle yields in  $D^0$ -tagged and inclusive jets as a function of  $\ln(1/\theta)$ , where  $\theta$  is the angle with respect to the jet axis, as measured in pp collisions at  $\sqrt{s} = 13$  TeV.

A detailed study of jet substructure has also been carried out in Pb-Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV [15] and has revealed that in nucleus-nucleus collisions jets are more collimated than in pp at the same energy, as seen in Fig. 6. This measurement indicates that the medium has a significant resolving power for splittings with a particular dependence on the coherence scale.

### 3. Emergence of QGP phenomena from QCD

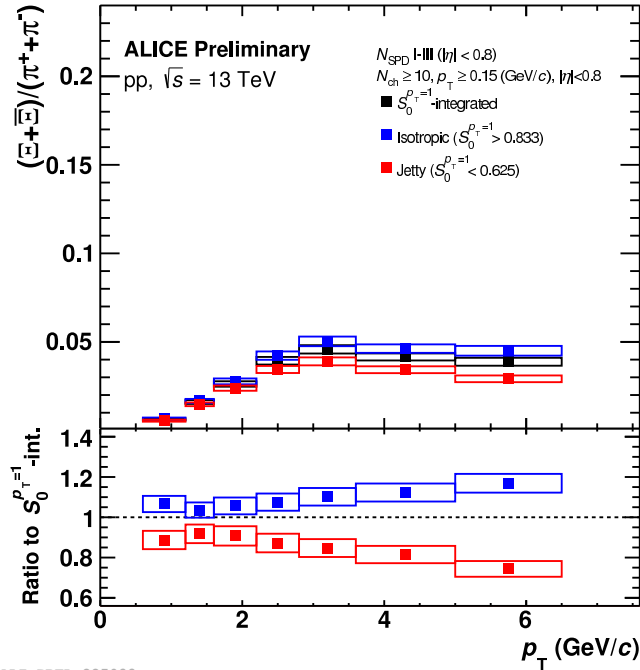
Substantial advances in our understanding of high density QGP were obtained in measurements performed in high-multiplicity pp and p-Pb collisions at the LHC. The observation of the double ridge [16–18] and strangeness enhancement [19] in such conditions were hallmark discoveries and their implications for heavy-ion physics is still being assessed by the community at large. To deepen our understanding of strangeness enhancement in proton-proton collisions, a double-differential study in which the  $\Xi/\pi$  ratio is studied in low- and high-sphericity, high multiplicity pp collisions has been performed [20]. Figure 7 shows that more isotropic events, which present a larger sphericity, are also observed to have a larger  $\Xi/\pi$  ratio, indicating that these are the ones that drive strangeness production higher and further constraining theoretical models.

Studies similar to the ones carried out in the light-flavour sector are also underway for heavy-flavour hadrons such as the  $\Lambda_c$  and the  $D^0$  meson [21]. These have been studied as a function of charged-particle multiplicity density in pp collisions at  $\sqrt{s} = 13$  TeV, with the  $p_T$ -differential  $\Lambda_c/D^0$  ratio exhibiting an increase at intermediate  $p_T$  that is reminiscent of that observed for the  $\Lambda/K_S^0$ , as can be seen in Fig. 8. When the  $\Lambda_c/D^0$  in a specific momentum interval is plotted as a function of the charged-particle multiplicity density in pp, p-Pb and Pb-Pb, as seen in Fig. 9, the progression seems to be smooth as a function of multiplicity, again similarly to the  $\Lambda/K_S^0$  [22]. When compared to PYTHIA8 simulations, the  $\Lambda_c/D^0$  can only be qualitatively described when



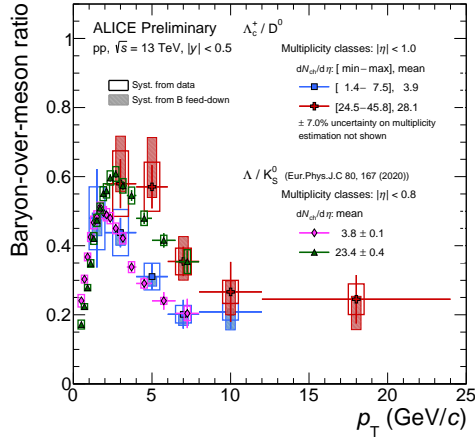
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**Figure 6:** Distribution of first jet splitting angle  $\theta_g$  measured in Pb-Pb collisions at  $\sqrt{s_{NN}}$  TeV and in pp collisions at the same energy using jet grooming techniques. Figure from [15].



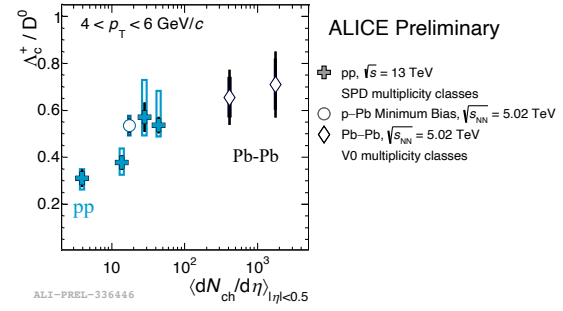
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**Figure 7:** Transverse-momentum-differential  $\Xi/\pi$  ratio in high- and low-sphericity pp collisions at  $\sqrt{s} = 13$  TeV.



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**Figure 8:** Transverse-momentum-differential  $\Lambda_c^+/D^0$  ratio as a function of charged-particle multiplicity density in pp collisions at  $\sqrt{s} = 13$  TeV shown together with the  $\Lambda/K_S^0$  ratio for comparison.



**Figure 9:** Multiplicity dependence of the  $\Lambda_c^+/D^0$  ratio in the momentum interval  $4.0 < p_T$  (GeV/c)  $< 6.0$  in pp collisions at  $\sqrt{s} = 13$  TeV and in p-Pb and Pb-Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV.

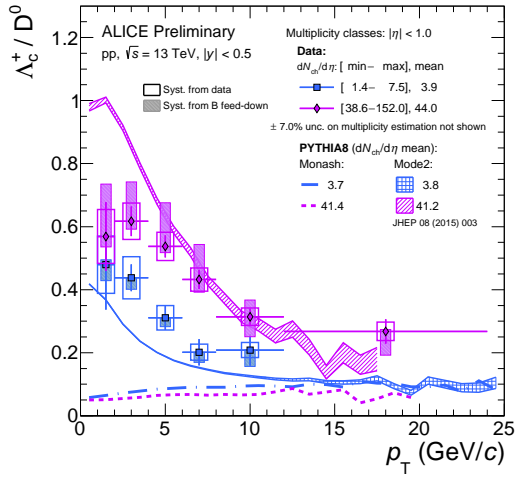
using color reconnection in ‘mode 2’, in which colour junctions are allowed to form [23], as can be seen in Fig. 10. Further measurements in Run 3 will allow for a significant reduction in both statistical and systematic uncertainties in this measurement [24].

Similar observations also hold for the  $p_T$ -differential  $\Sigma_c^{0,+}/D^0$  and  $\Xi_c^{0,+}/D^0$  [26], though for the latter even PYTHIA with junctions fails to reproduce the data by more than a factor of three. The  $\Xi_c^{0,+}/D^0$ , shown in Fig. 11, is also not correctly reproduced by a coalescence model [25]. These findings indicate that a lot is still to be understood regarding charmed baryon production in high-energy collisions, with more insights expected to come from the LHC Run 3 and beyond.

#### 4. Beyond heavy ions: The LHC as a general-purpose QCD laboratory

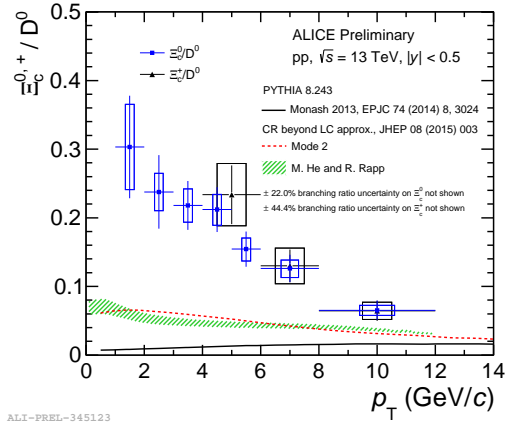
The excellent particle identification capabilities of the ALICE experiment, coupled with the copious production of hadrons at the LHC, leads to unique opportunities for advancing our understanding of QCD in general. One example of a study that leveraged these conditions is the first measurement of antideuteron and antihelium absorption [27], which led to cross section measurements at low momenta that are highly relevant to astrophysics. The preliminary result for the  ${}^3\overline{\text{He}}$  is shown in Fig. 12 and the corresponding result for antideuterons is published in [28]. Further advances in antinuclei measurements are expected in Run 3 and beyond.

Another very exciting opportunity that has presented itself at the LHC is to study the proton-hyperon interaction. This can be performed by measuring proton-hyperon pair momentum correlations, which, combined with a determination of the particle-emitting source size, allows for a precise comparison of the interaction characteristics to predictions from lattice QCD, as can be seen in Fig. 13. This measurement, published in Nature [29], further adds to our understanding of proton-hyperon interactions, which are relevant for constraining the equation of state of neutron



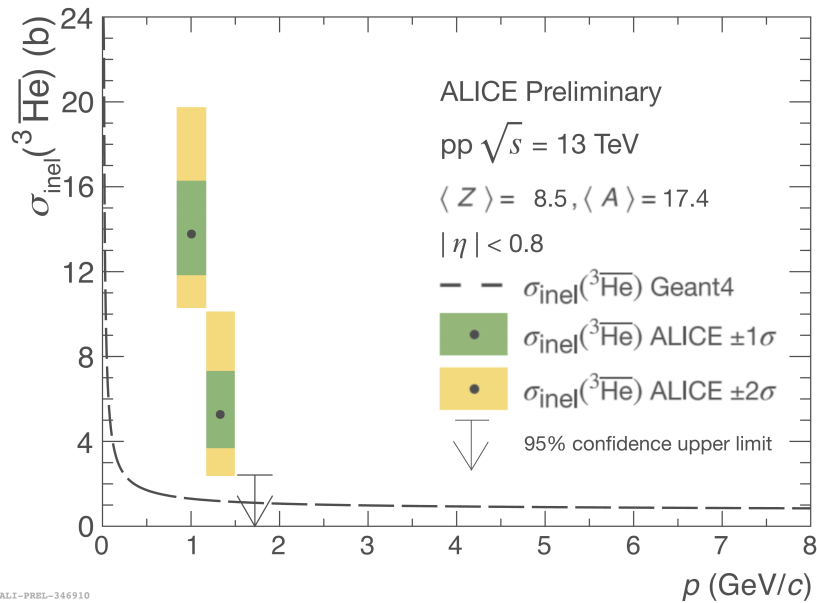
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**Figure 10:** Transverse-momentum-differential  $\Lambda_c^+/D^0$  ratio in two selected charged-particle multiplicity intervals in pp collisions at  $\sqrt{s} = 13$  TeV compared to predictions from PYTHIA8 using various settings of colour reconnection.



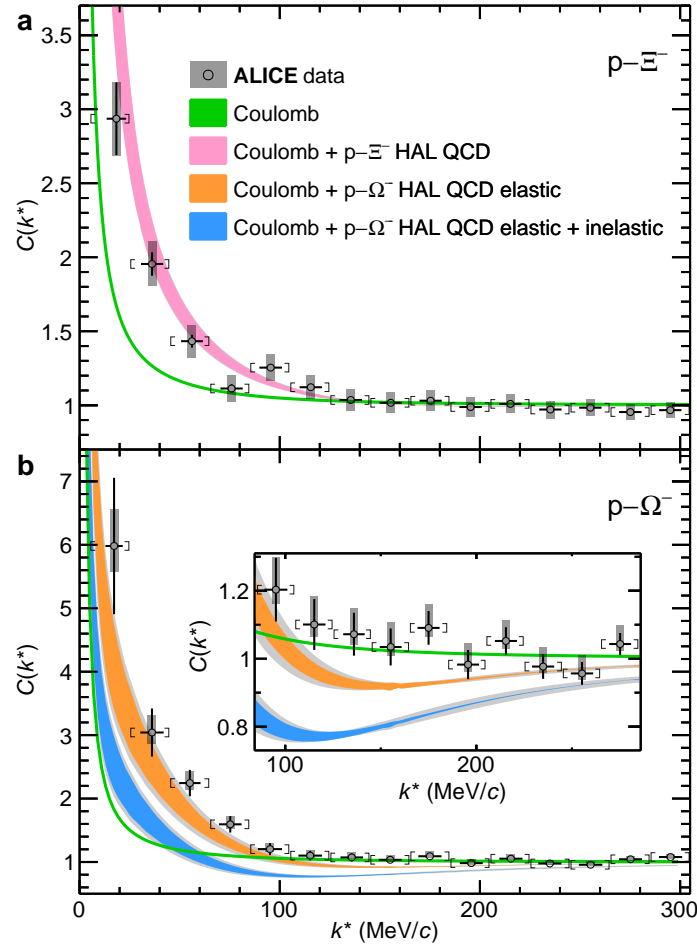
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**Figure 11:** Transverse-momentum-differential  $\Xi_c^{0,+}/D^0$  ratio in two selected charged-particle multiplicity intervals in pp collisions at  $\sqrt{s} = 13$  TeV compared to predictions from PYTHIA8 using various settings of colour reconnection as well as with predictions from the coalescence model from M. He and R. Rapp [25].



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**Figure 12:** Inelastic cross section  $\sigma_{\text{inel}}(^3\text{He})$  as a function of total momentum measured via secondary interactions with the material in pp collisions at  $\sqrt{s} = 13$  TeV.



**Figure 13:** Experimental  $p\text{-}\Xi^-$  and  $p\text{-}\Omega^-$  correlation functions measured in high-multiplicity pp collisions at  $\sqrt{s} = 13$  TeV compared to predictions from lattice QCD using various combinations of interaction terms. Figure from [29].

stars, and constitutes the first measurement of an attractive strong interaction between protons and  $\Omega$  baryons.

## 5. The bright future ahead

The ALICE collaboration has a comprehensive upgrade programme foreseen in the next years. During Long Shutdown 2 (LS2), the TPC readout is being replaced in its entirety by a GEM-based system that is able to cope with a factor 50 higher data rate. Furthermore, the inner tracking system (ITS), previously comprised of pixel, strip and drift detectors, is being replaced by a fully pixel-based system, effectively multiplying the number of channels by a factor of a thousand. The much improved ITS, coupled with a pixel-based muon forward tracker (MFT), will improve impact parameter resolution by a factor of 3 in Run 3. This detector setup will lead to significant advances in our understanding of heavy-flavour meson and baryon production, as well as enable studies of



(anti)(hyper)nuclei up to a mass number of 4. In addition, hadronic interactions up to  $\Omega - \Omega$  pairs will be within reach [30], and a first measurement of the low-mass dielectron spectrum will allow for an estimate of the QGP temperature before hadronization.

Further into the future, the Long Shutdown 3 will also bring about the installation of an even more advanced inner tracking system, the ITS3 [31], comprised solely of thin bent sheets of silicon in the active detection area. This groundbreaking detector will enable even higher precision studies of low- $p_T$  charm production, including the possibility of searching for charmed nuclei. In addition, a forward calorimeter, FoCal, is planned, and will allow for high-precision measurements of the gluon nPDF down to  $x \approx 10^{-5}$  and below [32].

Even beyond, a project for a next-generation heavy-ion experiment at the LHC, named ALICE 3, is being actively pursued [33]. It is to be a fast, ultra-thin detector solely based on silicon sensors, allowing for yet another fifty-fold increase in data rates. In so doing, ALICE 3 will be able to exploit higher LHC luminosities using nuclei lighter than lead, while still retaining particle identification capabilities via state-of-the-art silicon-based time of flight determination with  $\sim 20$  ps time resolution. This brand new detector will enable ultimate performance for multi-heavy-flavour hadron detection, thermal radiation via virtual photons from the QGP and soft hadrons (below  $p_T$  of 50 MeV/ $c^2$ ). This initiative has been supported in the latest European Strategy for Particle Physics Update, ensuring that the future of heavy-ion physics will remain strong even beyond the next decade.

## Acknowledgements

The author would like to acknowledge support from FAPESP via project number 17/05685-2 and CNPq via grant number 312661/2019-3.

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