

The Two Hemispheres Method for Multijet BSM Searches

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A new method for identifying hints of possible beyond the standard model (BSM) signals produced at the Large Hadron Collider (LHC) with high jet multiplicity final states is proposed. In particular, the QCD background is estimated in a data driven way. Based on the simplified picture where QCD multijet events are created from a $2 \rightarrow 2$ process followed by cascade branching of the outcoming partons, the proposed "Two Hemisphere Method" (THm) divides events to two hemispheres and predicts the distribution of the number of jets in a predefined high multiplicity signal region. Validation of the above-mentioned assumption was performed using LO, NLO, and NNLO simulations, showing no effect of higher order calculations on the prediction accuracy.

The sensitivity of the method was examined on topologically distinct scenarios of BSM multijet signatures and was able to show comparable sensitivity to other methods used in previous analyses. Since the sources of the uncertainties in this new approach are very different from the current methods, the procedures complement one another.

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Figure 1: Illustrations of a multijet event produced by (a) two back-to-back out-coming partons each followed by a parton shower (b) micro black hole decay (c) RPV SUSY gluino cascade decay chain.

1. Introduction

Several BSM models (such as Micro-Black-Holes (mBH) or R-Parity Violation (RPV) Supersymmetry) predict the possible production at the LHC of events with a large number of outcoming high energy partons. These events will give rise to final states consisting of a high multiplicity of energetic jets, i.e. multi-jet events. The identification of this type of signal through the observation of an excess of multijet events is nontrivial due to the presence of large SM background originating from QCD processes. Background estimation via simulation is limited and therefore a data driven approach is favorable.

2. Concept and Method

As a first order approximation QCD events can be seen as beginning with a $2 \rightarrow 2$ process that gives rise to two back-to-back out-coming partons, followed by a parton shower (Fig. 1a). Since the showering is carried out for each parton independently it is claimed here that at first approximation the number of jets in each hemisphere are not correlated. This hypothesis was validated for higher orders using NLO and NNLO generators in [1].

On the other hand, multijet signatures originating from massive BSM states (Fig. 1b and 1c) are expected to show stronger correlation between the jet multiplicities of each hemisphere. It is this difference that is utilized by the THm to predict the QCD background and differentiate a possible signal from that background. The procedure is outlined as follows:

Divide events into 2 hemispheres (A and B) using the plane perpendicular to the transverse thrust axis. The number of jets in each hemisphere will be denoted N_{Jets}^A and N_{Jets}^B respectively. For each value of N_{Jets}^A , obtain the distribution of N_{Jets}^B (illustrated in Fig. 2).

Under the simplistic $2 \rightarrow 2$ assupption the showering initiated by each outcoming parton is independant of the other outcoming parton, therefore the normalized N_{Jets}^{B} distributions should be similar for any chosen N_{Jets}^{A} . This is shown in Fig. 3 using a LO generator calculation integrated with a parton shower (Pythia) and was validated using NLO and NNLO generators in [1].

The hypothetical High N_{Jet} signal is unlikely to give rise to highly asymmetric events with a large jet multiplicity in hemisphere B and only one jet in hemisphere A. Therefore, the $N_{Jets}^{B}(N_{Jets}^{A}=1)$ distribution (black) should always be relatively signal free. $N_{Jets}^{B}(N_{Jets}^{A}=1)$ can thus be used to represent the N_{Jets}^{B} distribution of pure QCD and serve as a QCD background estimation for those samples which might host signal events (i.e. $N_{Iets}^{A}=2,3,4...$). An excess of events with high jet



Figure 2: Representation for the distribution of the number of jets in hemisphere B (N_{Jets}^{B} , right hemisphere) for each value of N_{Jets}^{A} (left hemisphere).



Figure 3: Normalized distributions of N_{Jets}^B for each value of N_{Jets}^A . As can be seen all distributions are in close agreement.

multiplicity may be considered as a possible indication for the presence of a signal (Fig. 4 bottom right).

3. Simulation

QCD events were generated and showered with PYTHIA8 [2] at $\sqrt{s}=13$ TeV using the Monash 2013 tune [3]. Jets were reconstructed using the *anti-kt* algorithm [4] implemented in the FastJet 3.2.1 package [5] with a radius parameter value of R = 0.4. All jets were required to satisfy $p_T > 50$ GeV and $|\eta| < 2.8$. mBH signal samples were produced using a semi-classical approximation with the CHARYBDIS2 [6] event generator. The signal sample parameters include the number of extra dimensions (n), the value of the diminished Planck mass (M_D) and the threshold above which the semi-classical treatment of the microscopic BH is expected to be valid (M_{th}). The production crosssection for each sample is determined by its parameters. Multijet RPV SUSY signal samples were produced using the MadGraph5_aMC@NLO [7] event generator, parton showered with PYTHIA 8. Cascade gluino decay samples were produced for values of 1000 $< m_{\tilde{g}} < 2100$ GeV and neutralino mass 50 GeV $< m_{\tilde{\chi}_1^0} < 1.65$ TeV (where always $m_{\tilde{\chi}_1^0} < m_{\tilde{g}}$).

4. Results

In order to test the sensitivity of the above outlined search procedure, simulated signals of two topologically distinct scenarios, mBH and RPV SUSY, were added to the simulated QCD sample. These simulated signal events revealed that the mBH samples have lower average jet multiplicty and larger variance thus giving rise to a non-negligible signal contamination in the supposedly signal free N_{Jets}^{A} = 1 (Fig. 4 top left). Such signal contamination of the conjectured "signal-free" prediction deems that the THm sensitivity for these mBH samples is much reduced (Fig. 4 bottom left). The RPV SUSY signal samples were found to have higher average jet multiplicity and a smaller variance thus leading to negligible contamination in N_{Jets}^{A} = 1 (Fig. 4 top right) and therefore provide a signal that can potentially be detected by the THm (Fig. 4 bottom right). Figure 5 shows the expected limit at 95% CL using the Two Hemispheres method (blue) overlaid on the exclusion contour from the ATLAS RPV multijet analysis [8]. Exclusion contours given in the $(m_{\tilde{g}}, m_{\tilde{\chi}_{1}^{0}})$ plane for the



Figure 4: Top:N^A_{Jets}vs. N^B_{Jets}of characteristic mBH sample (left) showing significant (~20% of total sample) signal contamination in N^A_{Jets}=1 (red rectangle). RPV SUSY (right) showing negligible contamination. Bottom: Normalized N^B_{lets}distributions for SM bkg (Pythia) injected with signal. RPV SUSY (right) causes N_{Jets}^{A} = 3,4 (green, blue) to deviate significantly from the N_{Jets}=1, thus exposing signal as desired. mBH (left) suffers from significant signal contamination in N_{Jets}^A =1 thus greatly reducing sensitivity

Figure 5: Expected limits (95% CL) for the RPV SUSY gluino cascade decay model using the THm (blue) compared to those achieved by the ATLAS RPV multijet analysis (green).

gluino cascade decay model. The results are comparable to those achieved in the ATLAS paper (interpolation between the points of the signal parameter grid was not perfomed for the THm). The THm is new and further optimization may improve sensitivity. Yet, even with similar sensitivity the two analyses, being so different, complement each other. Should a signal be detected in one analysis the other can be used to confirm or refute its existence.

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