

## The effect study of thunderstorm electric field on the lateral distribution of positrons and electrons at YBJ

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*Abstract.* For a high altitude experiment, the lateral distribution (particle density distribution) of secondary particles at detector level is an important parameter to get the information of primary shower event. During thunderstorms, the charged particles in extensive air showers will be accelerated or decelerated in strong atmospheric electric field. The intensity and the position of ground cosmic rays could be influenced. In this work, Monte Carlo simulations are performed by using CORSIKA to study the effect of thunderstorms electric field on the lateral distribution of charged particles at YBJ (4300 m a.s.l., Tibet, China). We found that the lateral distribution will be changed significantly in thunderstorm field. And the variation amplitude is found to be highly dependent on the electric field, the zenith angle of the primary shower event and the distance to the shower core. Our simulation results could be helpful in understanding the phenomena observed in ground based experiments (such as ARGO-YBJ) and the physical mechanism of secondary charged particles accelerated by thunderstorm electric field.

*Keywords:* Thunderstorm electric field, Cosmic rays, Lateral distribution, Monte Carlo simulation

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## 1. Introduction

When a primary cosmic ray enters the atmosphere, numerous secondary particles, such as positrons, electrons, photons and muons, are created through the process of hadron and electromagnetic cascades. These secondary particles are distributed over a large area around the shower core. The lateral distribution, which is the density distribution of ground cosmic rays around the shower axis at detector level, is crucial to study the information of primary shower event, such as primary components, energy spectrum and the acceleration mechanism. Therefore, scientists have been paying close attention to the lateral distribution of secondary particles.

The Nishimura-Kamata-Greisen (NKG) function [1, 2] gives a good theoretical approach for the lateral distribution of electromagnetic particles, from which the dependency of the lateral growth of cascade on the energy, intensity and type of the primary cosmic rays can be clearly shown. Several ground-based observation experiments, such as ARGO-YBJ [3], KASCADE [4], NBU [5] and Utah neutrino detector [6], have reported the lateral distribution of secondary particles associated to the energy of primary cosmic rays. A few simulations have been done to study the lateral distribution. Apel et al. [4] found that the lateral distributions of electrons and muons was related to the mass, energy and type of the primary shower event. And the same laws were obtained by Zhao [7], Wang [8] and Atreidis [9].

In addition, the influence of the environment is not to be neglected. Sampath et al. [10] indicted that the geomagnetic field and atmospheric air density have considerable effects on the lateral distribution of extensive air shower. During thunderstorms, the magnitude of atmospheric electric field is commonly up to 1000 V/cm [11–13]. In such strong fields, the location, time, energy and flux of ground cosmic rays will be changed significantly. It means that the lateral distribution of secondary particles, which produced by primary cosmic rays interact with the atmosphere, will also be inevitably influenced. In this work, Monte Carlo simulations are performed by using CORSIKA to study the effect of near-earth thunderstorm electric field on the lateral distribution of positrons and electrons at YBJ. This paper is organized as follows: The simulation parameters are introduced in Section 2. Simulation results and discussions are shown in Section 3. The conclusions are presented in Section 4.

## 2. Simulation parameters

The CORSIKA (COsmic Ray SIMulations for KAScade) is a Monte Carlo simulation program to study the evolution and properties of EAS in the atmosphere initiated by nuclei, hadrons, photons, or any other particles [14, 15]. In this work, we use the 7.5700 version of CORSIKA. The selected hadronic interaction model is QGSJET-II-04 (Quark Gluon String model with JETs) for high energy and GHEISHA (Gamma Hadron Electron Interaction SHower code) in the low energy range. The energy cutoff was set to 0.1 MeV, below which value positrons and electrons are discarded from the simulation.

In this work, the electric fields are regarded as homogeneous and the thickness of the field layer is 1000 m, from the detector level (4300 m) up to 5300 m, corresponding to the atmospheric depth from 606 g/cm<sup>2</sup> to 531 g/cm<sup>2</sup>. We assume proton primary with a zenith angle in the interval from 0° to 60°. Number of showers is 2×10<sup>4</sup>. According to the observed energy range of the ARGO-YBJ detector, we simulated primary protons with energy from 100 GeV to 1 PeV and power law spectrum of -2.7. The electric fields are chosen as a series of values in the range of -1000 V/cm to 1000 V/cm. Here, we define the positive electric field as one that accelerates positrons downward in the direction of the earth.

### 3. Simulation results and discussions

In this paper, the effects of thunderstorm electric fields on the lateral distributions of positrons and electrons at YBJ are simulated.

#### 3.1 The lateral distributions of positrons and electrons in absence of electric field

Fig. 1 shows the lateral distributions of positrons and electrons in absence of electric field at YBJ. We can see that the percent of particle density decreases as the  $R$  (the radius of secondary particles to shower core) increases. The average radius can be calculated by

$$\bar{R} = \sum_{i=1}^N R_i / N \quad (3.1)$$

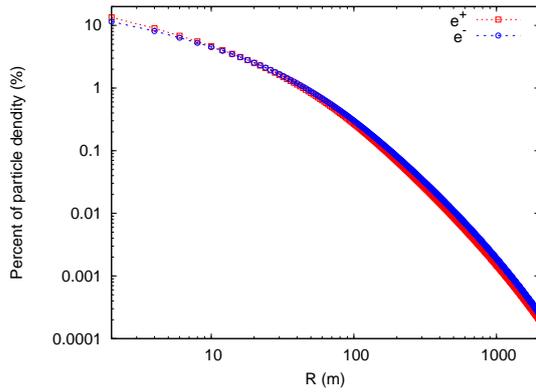
here  $R_i$  is the distance of the particle  $i$  to shower core at detector level,  $N$  is the total number of particles in the shower. In the following figures, the distributions are normalized to the density at  $R = 2000$  m. From Fig.1, we can obtain the average radius  $\bar{R}$  of positrons and electrons are 321.1 m and 340.9 m, respectively. That means the lateral distribution of electrons is wider than that of positrons.

Fig. 2 shows the density distribution of particles as a function of the distance to shower core with different primary zenith angles in absence of electric field. We can see the lateral widens with the higher zenith angle. The average radius are 282.1 m ( $\theta \sim 0^\circ - 20^\circ$ ), 323.7 m ( $\theta \sim 20^\circ - 40^\circ$ ) and 457.5 m ( $\theta \sim 40^\circ - 60^\circ$ ), respectively.

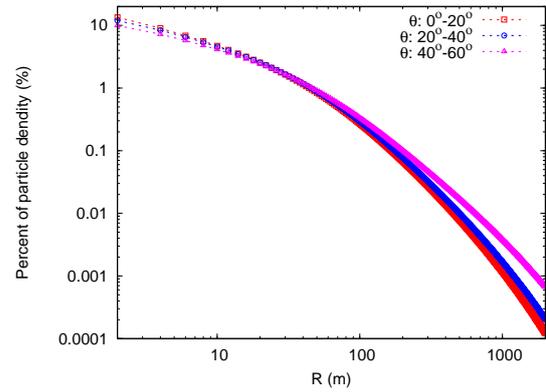
#### 3.2 The variations of lateral distributions in thunderstorm fields

In the field strength of 1000 V/cm, the distributions of particle density as a function of the distance to shower core are shown in Fig. 3. The result in absence field is plotted just for comparison with the plots in fields. It is easy to see that the lateral distributions become wider in fields. In 1000 V/cm, the average lateral distribution is 348.1 m, and the increase amplitude is about 4.4%. In -1000 V/cm, the average radius is 364.5 m, and the amplitude enhancement is up to 9.3%.

For different primary zenith angles, the distribution of particle density in 1000 V/cm and -1000 V/cm are shown in Fig. 4 and Fig. 5, respectively. As shown in the two figures, the effects of

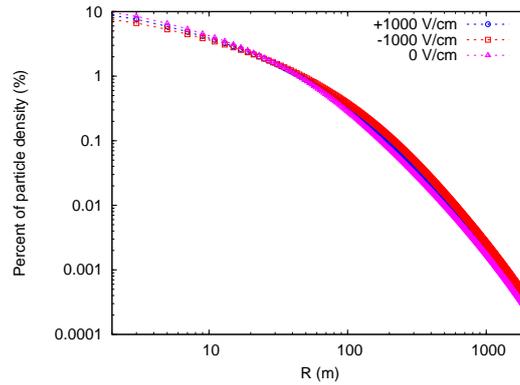


**Fig. 1:** Lateral distributions as a function of distance to shower core.



**Fig. 2:** Lateral distributions as a function of distance to shower core with different primary zenith angles.

electric field on the lateral distribution with different primary zenith angles are different. And the effect becomes larger as the primary zenith angle increases. The average radius in different fields with different primary zenith angles are listed in table 1.

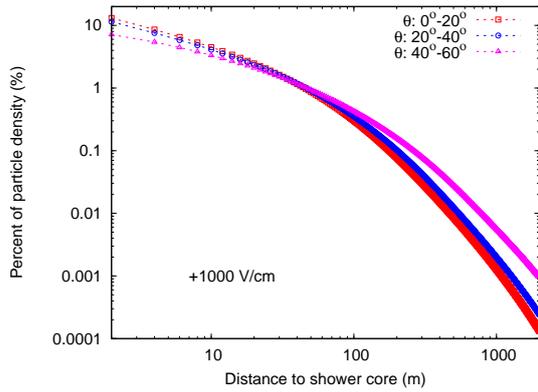


**Fig. 3:** Lateral distributions as a function of distance to shower core in  $\pm 1000$  V/cm.

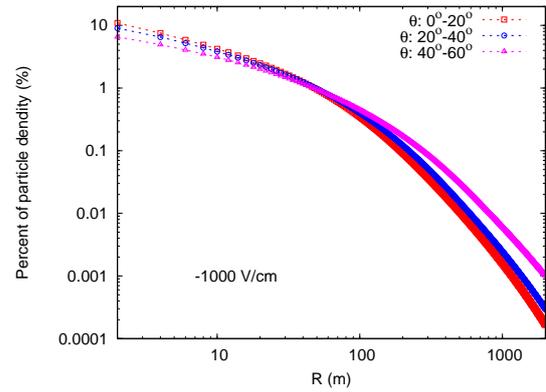
**Table 1:** the  $\bar{R}$  in different fields with different primary zenith angles

$E$ (V/cm) \ $\bar{R}$ (m)	$\theta: 0^\circ - 20^\circ$	$\theta: 20^\circ - 40^\circ$	$\theta: 40^\circ - 60^\circ$
0	282.1	323.7	457.5
+1000	286.9	337.1	489.0
-1000	301.0	353.6	502.1

Fig. 6 shows the average lateral distributions in different electric fields at YBJ and Fig. 7 shows the variations of lateral distributions. We can see that the average radius increases as the electric field strength increases. The average radius is larger in negative field than that of in the same strength positive field.

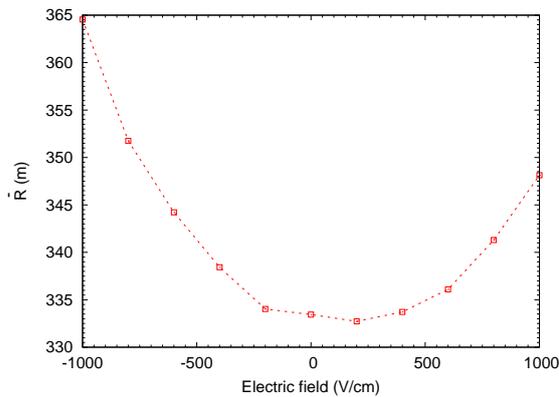


**Fig. 4:** Lateral distributions as a function of distance to shower core with different primary zenith angles in +1000 V/cm.

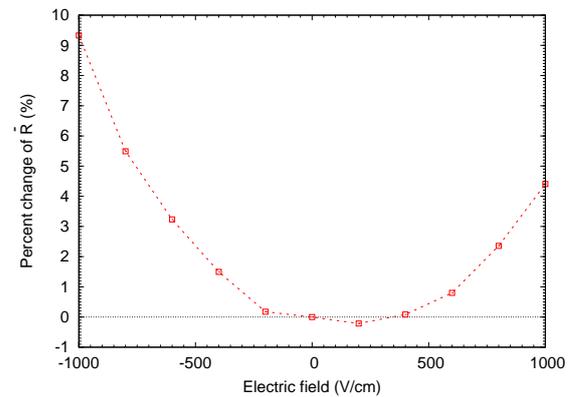


**Fig. 5:** Lateral distributions as a function of distance to shower core with different primary zenith angles in -1000 V/cm.

Fig. 8 and Fig. 9 show the average lateral distributions as a function of electric field with different primary zenith angles. The average radius gradually increases with increasing field strength. In addition, for the larger zenith angle, the average radius become much larger. That means the enhancement amplitude is higher for showers with larger zenith angles.



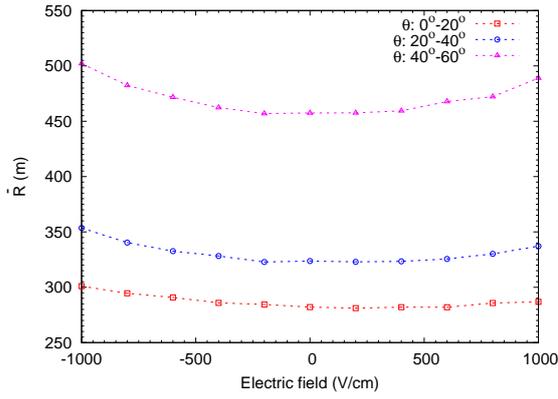
**Fig. 6:** The distributions of average radius as a function of field.



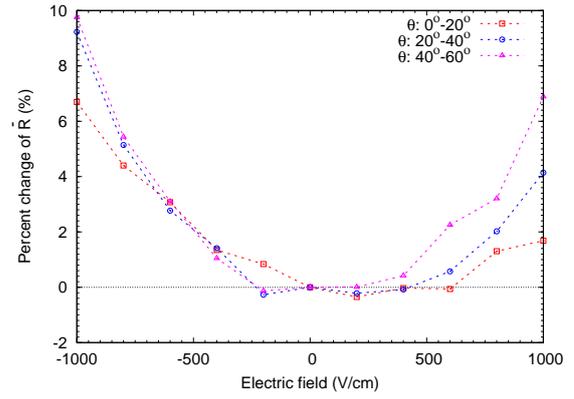
**Fig. 7:** Percent change of the average radius as a function of field.

### 3.3 The variations of counting rate in electric field

Fig. 10 shows the variations of particle number as a function of the distance to shower core in  $\pm 1000$  V/cm. We can see the counting rate decrease in negative fields when  $R$  is less than  $\sim 10.6$  m, otherwise the number increase as the distance to shower core increases. There is a same law in positive field, and the particle number decrease if  $R$  is less than  $\sim 50.5$  m. The amplitude enhancement in a negative field is much larger than that of in positive field with the same strength, and the decrease amplitude is larger in positive field.

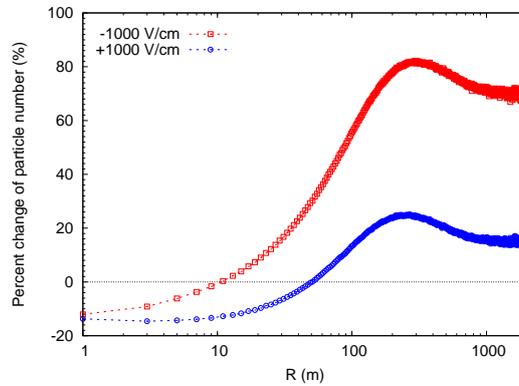


**Fig. 8:** The distributions of average radius as a function of field with different primary zenith angles.



**Fig. 9:** Percent change of the average radius as a function of field with different primary zenith angles.

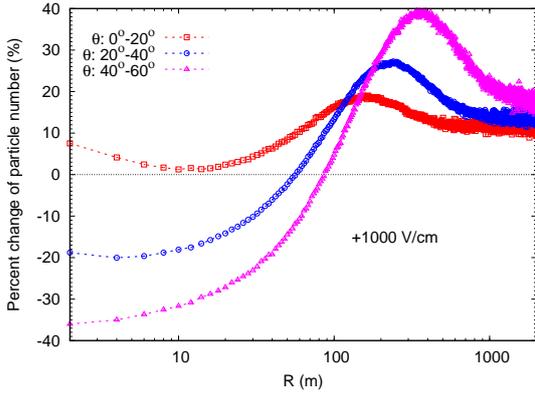
Fig. 11 and Fig. 12 show the percent change of particle number for different primary zenith angles in 1000 V/cm and -1000 V/cm, respectively. We can see the increasing phenomenon appeared in the area far from the shower core, and the amplitude enhancement becomes larger with the higher zenith angle. However, the decreasing phenomenon occurred in lower lateral radius range with higher zenith angle, and the declining amplitude increases with the increasing zenith angle.



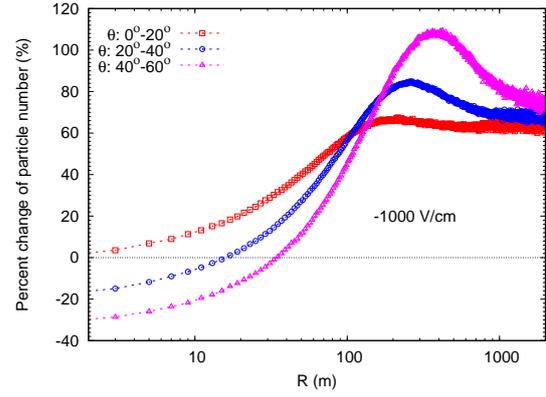
**Fig. 10:** Percent change of particle number as a function of distance to shower core in  $\pm 1000$  V/cm.

From our simulations, we can see the variations of particles number are highly dependent on the detector size. Fig. 13 describes the variations of particle number with different primary zenith angles in different fields at  $R \leq 60$  m. The counting rate increases in negative fields and decreases in positive ones, and the amplitudes depend on the zenith angle. The fluctuations in our results are due to the limited shower number in the Monte Carlo simulation.

At  $R \leq 2000$  m, the number variations as a function of electric field with different primary zenith angles are shown in Fig. 14. In negative electric fields and in positive fields greater than 600 V/cm, the number increases with increasing field strength, and the amplitude increases with

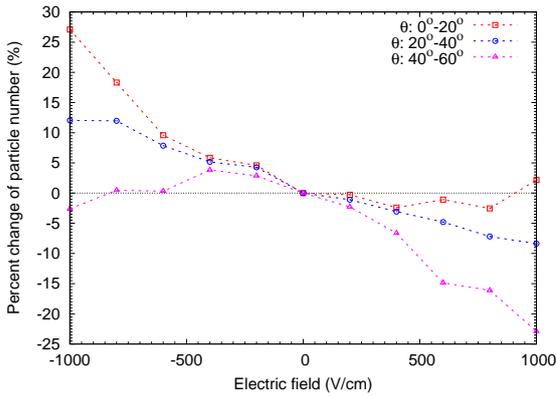


**Fig. 11:** Percent change of particle number as a function of distance to shower core with different primary zenith angles in +1000 V/cm.

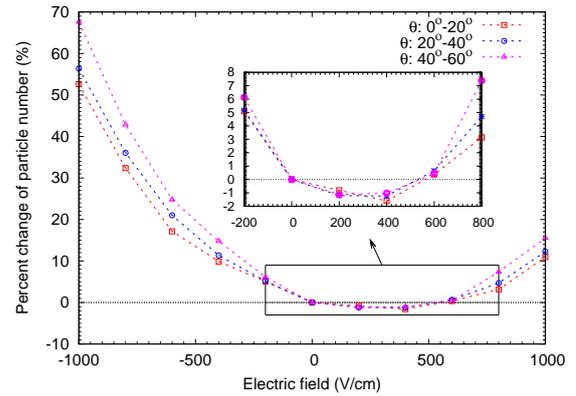


**Fig. 12:** Percent change of particle number as a function of distance to shower core with different primary zenith angles in -1000 V/cm.

the higher zenith angle. The amplitude of the enhancement in negative fields is significantly higher than that in positive ones with the same strength. The decreasing phenomenon occurs in the field range of 0 – 600 V/cm and the maximum amplitude is about 1.52%.



**Fig. 13:** Percent change of particle number as a function of electric field with different primary zenith angles at  $R \leq 60$  m.



**Fig. 14:** Percent change of particle number as a function of electric field with different primary zenith angles at  $R \leq 2000$  m.

#### 4. Conclusions

In this paper, Monte Carlo simulations are performed with CORSIKA 7.5700 packages to study the effect of near-earth thunderstorm electric field on the lateral distribution of positrons and electrons at YBJ. The simulation results show the lateral distribution widens in field, and the amplitude increases with the increasing strength of the field and the primary zenith angle.

The counting rate variations are highly dependent on the polarity of electric field, the zenith angle of the primary shower event, and the detector size (the distance to shower core). In negative

field, if the  $R$  becomes large enough, the counting rate increases and the amplitude enhancement is larger for higher zenith angle; otherwise, the amplitude enhancement is larger for lower zenith angle. In positive field, if  $R$  is less than a certain value, the decreasing phenomenon occurs, and the amplitude becomes larger for higher zenith angle.

The simulation results of field effects on the lateral distributions will be useful in understanding the phenomena observed in ground based experiments.

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