

A First Look at Periodicity in HAWC with TeV Binaries

Chad Brisbois, Petra Huentemeyer, Henrike Fleischhack, Binita Hona

Michigan Technological University

Email: cbrisboi@mtu.edu, petra@mtu.edu, hfleisch@mtu.edu, bhona@mtu.edu

Chang Rho

University of Rochester

E-mail: crho2@ur.rochester.edu

for the HAWC collaboration ‡

Only five binary systems have been found to emit at TeV energies. Each of these systems is composed of a massive O or B type star and a compact object (black hole or a pulsar). The type of compact object and the origin of the gamma-ray emission is unknown for most of these systems. Extending spectral observations to higher energies can help disentangle the nature of the compact object as well as the particle acceleration mechanisms present. Interestingly, the TeV emission from these systems does not always coincide with their emission in GeV or X-ray, which is how many such systems have been originally discovered. Increased coverage of these systems may allow HAWC to see precisely when in the orbit the TeV emission begins and ends. The HAWC Observatory detects TeV gamma-rays with high sensitivity, covering over two-thirds of the overhead sky every day. Applying a stacking method to known TeV binary systems can help HAWC enhance the signal from TeV binaries above the steady background from other sources in the galaxy. We will present results from this stacking analysis using 760 days of HAWC data.

35th International Cosmic Ray Conference — ICRC2017 10–20 July, 2017 Bexco, Busan, Korea

^{*}Corresponding Author

[†]Speaker.

[‡]See http://www.hawc-observatory.org/collaboration/icrc2017.php for full author list.

1. Introduction

Binary systems which emit γ -rays at TeV energies are exceedingly rare objects. Only a handful are known, and are listed in Table 1. It is presently unknown whether most of these objects are driven by an accreting black hole (also called the microquasar scenario) or a pulsar wind nebula. PSR B1259-63 represents a notable exception to this trend, because it is known that the compact companion is a pulsar. The search for more of these systems is a primary focus of discovery, because both of these cases are likely to be found. It is possible that many massive binary systems go through a γ -ray emitting phase, but due to their relatively short life as a TeV binary ($\sim 10^5$ years) there are only a few in our galaxy that are observable at any given time [1].

Name	Period (days)	Distance (kpc)
LS 5039 [2]	3.9	2.5
HESS J0632+057 [3]	315	1.6
LS I +61 303 [4]	26.5*	1.11
HESS J1018-589 A [5]	16.58	5
PSR B1259-63 [2]	1241	2.3

Table 1: γ -ray Binaries listed in TeVCat [http://tevcat.uchicago.edu/]. Period and distance data are quoted from the references specified adjacent to their name. * LS I +61 303 has two periodicities, a further discussion of which can be found in [4]. Of these sources, only LS 5039 is within HAWCs FoV.

The High Altitude Water Cherenkov observatory (HAWC) is located on the Sierra Negra Volcano in the Pico de Orizaba National Park in Mexico. Due to its latitude (+19°N), elevation (4100m), and detection method (water cherenkov detection) it is sensitive to TeV emission on an extremely high duty cycle across 2/3 of the sky each day [6]. This high uptime (>90%) allows HAWC to be sensitive to time variability in emission that other instruments cannot observe. HAWC has a real-time flare monitoring system which is capable of detecting flares on time scales as small as two minutes [8]. These qualities make HAWC an excellent instrument for observing time-dependent behavior in the TeV sky. Presently HAWC uses what is referred to as fractional nHit bin $(0 \le \mathcal{B} \le 9)$ analysis [6], where energy proxy bins are defined by the fraction of the detector triggered by a particular air shower.

LS 5039 is a high-mass X-ray binary system composed of a massive O-type star (M = $23M_{\odot}$) and a compact object (M = $3.7M_{\odot}$) with an orbital period of 3.9 days [2]. The nature of the compact companion is unknown. A variety of instruments such as H.E.S.S. and Fermi-LAT have detected γ -ray emission from LS 5039. In [9], H.E.S.S. was able to observe two distinct spectral shapes in the γ -ray emission of LS 5039 depending on the orbital phase (0 < ϕ ≤ 1) of the system. In the high state (0.45 < ϕ ≤ 0.9), LS 5039 exhibited a hard power law spectrum with a cutoff, and in the low state (0.45 ≤ ϕ & ϕ > 0.9) a simple power law with a softer index. A plot of the spectral energy distribution can be found in Figure 1 which was taken from [9]. The same definitions of high and low state are used in the present analysis.

From Figure 1 it is ambiguous what the spectrum beyond 10 TeV looks like. If the high state truly cuts off beyond this energy range, or if the high state emission is simply "on top of" the low state emission is an open question that HAWC may be able to address. HAWC is currently the only instrument that has hundreds of days of exposure time on sources at energies greater than 10 TeV.

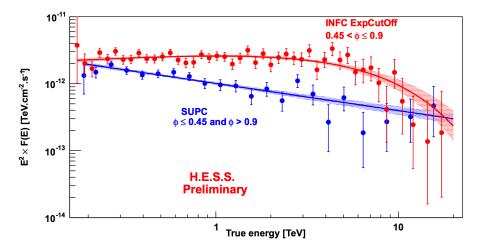


Figure 1: Spectrum of LS 5039 binned into high and low state, taken from [2].

LS 5039 is embedded in a busy region along the galactic plane. There are two bright extended sources nearby, HESS J1825-137 and HESS J1826-130. HAWC detects a significant γ -ray excess over the expected background at the location of LS 5039 (see Figure 2). However, it is not clear how much of this excess is due to γ -ray emission from LS 5039 and how much is due to leakage from the other two sources. In this presentation, we will explore a method to measure the difference in flux between the high and low state of LS 5039 as seen by HAWC.

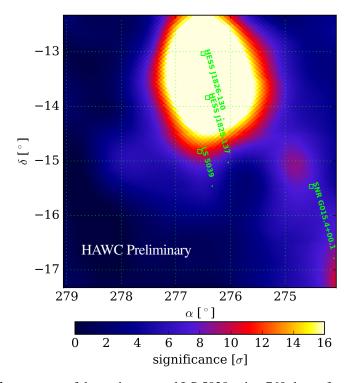


Figure 2: Significance map of the region around LS 5039 using 760 days of accumulated HAWC data.

2. Method

HAWC events are binned according to the phase of LS 5039. The data is analyzed with the standard HAWC analysis software [6] [7]. The two flux maps are produced by fitting a single point source with a power-law spectrum at the position of each pixel, with the spectral index fixed to -2.0. The quoted fluxes refer to the fitted differential flux at the pivot energy of 7 TeV. The high state map contains 342 days of data, and the low state contains 418 days.

By further subdividing the two sets of events for the high and low states of the binary into bins corresponding to the low-energy nHit bins and the high-energy nHit bins, we can gain a small window into the relative flux between the two states. If there is a positive relative flux in the low energy bins, and a negative relative flux in the high energy bins, then that is consistent with the high state flux not simply being "on top of" the low state flux. If at high energies the difference is zero, then it possible that the process generating the high state turns off at those energies, but the low state process persists.

Assuming the uncertainty on the flux normalization is Gaussian and purely statistical, the significance of the flux difference between the high and the low state maps is estimated by Equation 2.1. The difference between the flux in the high (F_{high}) and $low(F_{low})$ states gives an estimate of the significance by dividing by the uncertainty (σ) in the measured flux at high and low energies.

$$\mathfrak{S} = \frac{F_{high} - F_{low}}{\sqrt{\sigma_{high}^2 + \sigma_{low}^2}} \tag{2.1}$$

This procedure is tested by applying it to a non-periodic source, the Crab nebula. Applying Equation 2.1 to the high and low state maps, the Crab nebula has disappeared from the maps, see Figure 3. The result of this important check is to be expected [6], considering that the periodicity is fixed to the periodicity of LS 5039. In Figure 3b, we see that the distribution is indeed approximately Gaussian, and therefore $\mathfrak S$ is treated as significance.

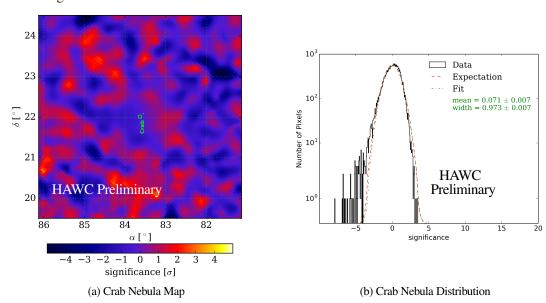


Figure 3: Significance of flux differences in Crab Nebula Region.

3. Results

In the region around LS 5039, the high and the low state maps display structure that is similar to what is seen in the full data set (see Figure 2), with nearby HESS sources dominating the emission of the region. The morphology of HESS J1825-137 and HESS J1826-130 are not known precisely enough to disentangle their emission from any eventual signal from LS 5039. Therefore we attempt to distinguish the relative flux between the two states, doing a per pixel subtraction of the flux according to equation 2.1. Following this prescription, the high and low state maps, can be converted to the relative flux between the states, as seen in Figure 4. The fluxes at the location of LS 5039 are listed in Table 2. No significant detection of a difference between the high and low state of LS 5039 is observed, nor are there significant detections elsewhere in the region.

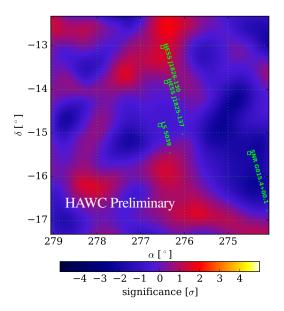


Figure 4: Overall relative flux between high and low state maps.

Мар	Flux $(10^{-14} TeV^{-1}cm^{-2}s^{-1})$
High State	1.13 ± 0.25
Low State	1.32 ± 0.24
Difference	-0.19 ± 0.35

Table 2: Differential flux at 7 TeV measured at the location of LS 5039 in the 'high' and 'low' states as well as the flux difference between the two states. Spectral fit was performed over the full energy range. Statistical uncertainties only. The quoted flux values have not been corrected for leakage from HESS J1825-137 and HESS J1826-130 and should not be taken as a measurement of the flux from LS 5039.

Splitting into the high and low energy maps for the relative flux between the states, corresponding to the crossover point near 10 TeV, there is again no significant difference between the two flux states (See Figure 5). All variations are consistent with statistical fluctuations. The flux measurements at LS 5039 are

summarized in Table 3. One critical caveat to the flux measurements listed in Tables 2 & 3 is that they contain any emission from LS 5039 as well as contamination from other sources.

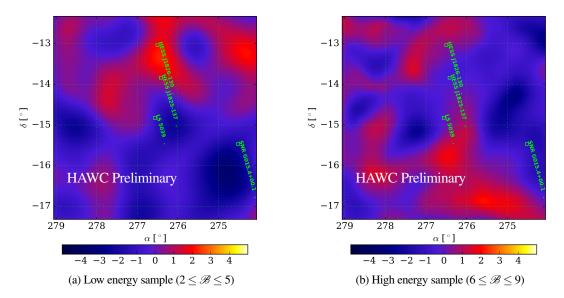


Figure 5: Significance of flux differences in high and low state maps in the region around LS 5039

Мар	Flux Difference $(10^{-15} \text{ TeV}^{-1} \text{cm}^{-2} \text{s}^{-1})$	8
High nHit bins $(6 \le \mathcal{B} \le 9)$	-0.59 ± 3.2	-0.18
Low nHit bins $(2 \le \mathcal{B} \le 5)$	-6.2 ± 8.4	-0.75
All nHit bins $(2 \le \mathcal{B} \le 9)$	-2.0 ± 3.5	-0.54

Table 3: Flux differences at the location of LS 5039, along with significance (\mathfrak{S}) for the full energy range as well as $2 < \mathcal{B} < 5$ and $6 < \mathcal{B} < 9$.

4. Discussion

We have presented a method to measure the flux differences between the high and low states of LS 5039. The method was validated by applying it to the crab nebula, a strong γ -ray sources that does not exhibit periodic variability. In Figures 4 & 5, the relative flux was consistent with zero in every case. This analysis does not detect any significant difference in the flux between the high and low states for LS 5039. However, the systematics of the method itself are being investigated, with the intention of producing upper limits in the future. As HAWC transitions from a fractional nHit bin based analysis to an analysis using energy estimators for measuring the spectrum, the ability to make precise statements about the flux (and variation thereof) at the location of LS 5039 or other variable sources for different energy bins will make HAWC more sensitive to differences in the differential flux. Additionally, efforts to measure the γ -ray emission from LS 5039 by modeling the morphology of the two neighboring sources

are currently underway. Planned improvements in HAWC's angular resolution due to improvements in shower reconstruction as well as the planned outrigger array will also help disentangle the emission in this complex region.

5. Acknowledgements

We acknowledge the support from: the US National Science Foundation (NSF); the US Department of Energy Office of High-Energy Physics; the Laboratory Directed Research and Development (LDRD) program of Los Alamos National Laboratory; Consejo Nacional de Ciencia y Tecnología (CONACyT), México (grants 271051, 232656, 260378, 179588, 239762, 254964, 271737, 258865, 243290, 132197), Laboratorio Nacional HAWC de rayos gamma; L'OREAL Fellowship for Women in Science 2014; Red HAWC, México; DGAPA-UNAM (grants RG100414, IN111315, IN111716-3, IA102715, 109916, IA102917); VIEP-BUAP; PIFI 2012, 2013, PROFOCIE 2014, 2015; the University of Wisconsin Alumni Research Foundation; the Institute of Geophysics, Planetary Physics, and Signatures at Los Alamos National Laboratory; Polish Science Centre grant DEC-2014/13/B/ST9/945; Coordinación de la Investigación Científica de la Universidad Michoacana. Thanks to Luciano Díaz and Eduardo Murrieta for technical support. C. Brisbois acknowledges support under the U.S. Department of Energy, Office of Science, Office of Workforce Development for Teachers and Scientists, Office of Science Graduate Student Research (SCGSR) program.

References

- [1] Dubus, G. 2013, Astronomy & Astrophysics Reviews, 21, 64
- [2] P. Bordas *et al.* [H.E.S.S. Collaboration], AIP Conf. Proc. **1792**, no. 1, 040017 (2017) doi:10.1063/1.4968921 [arXiv:1610.03264 [astro-ph.HE]].
- [3] S. Schlenstedt [VERITAS Collaboration], AIP Conf. Proc. **1792**, no. 1, 040023 (2017) doi:10.1063/1.4968927 [arXiv:1610.03751 [astro-ph.HE]].
- [4] Massi, M., Jaron, F., & Hovatta, T. 2015, Astronomy & Astrophysics, 575, L9
- [5] H. E. S. S. Collaboration, Abramowski, A., Aharonian, F., *et al.* 2015, Astronomy & Astrophysics, 577, A131
- [6] A. U. Abeysekara et al. 2017, The Astrophysical Journal, 843, 39
- [7] A. U. Abeysekara et al. 2017, The Astrophysical Journal, 843, 40
- [8] Abeysekara, A. U., Alfaro, R., Alvarez, C., et al. 2017, arXiv:1704.07411, in print
- [9] C. Mariaud *et al.* [H.E.S.S. Collaboration], PoS ICRC **2015**, 885 (2016) [arXiv:1509.05791 [astro-ph.HE]].