

A DIS Event Shape at N³LL *

Daekyoung KANG[†]

Theoretical Division, MS B283, Los Alamos National Laboratory, Los Alamos, NM 87545, USA

E-mail: kang1@lanl.gov

Christopher Lee

Theoretical Division, MS B283, Los Alamos National Laboratory, Los Alamos, NM 87545, USA

E-mail: clee@lanl.gov

Iain W. Stewart

Center for Theoretical Physics, Massachusetts Institute of Technology, Cambridge, MA 02139, USA

E-mail: iains@mit.edu

A high precision calculation of the event shape DIS thrust, with next-to-next-to-next-to-leading-logarithmic resummation and a rigorous treatment of hadronization corrections, is presented. Perturbative resummation uncertainties in the cross section are reduced to the 2% level for a significant region of the HERA phase space in x and Q , thus allowing for new accurate measurements of $\alpha_s(m_Z)$.

XXIII International Workshop on Deep-Inelastic Scattering

27 April - May 1 2015

Dallas, Texas

*LA-UR-15-25377, MIT-CTP 4640

[†]Speaker.

Event shapes provide a key method of measuring jets in deep-inelastic scattering (DIS). This was done successfully by H1 and ZEUS [1, 2, 3, 4, 5, 6] and compared with theoretical calculations with next-to-leading-logarithmic (NLL) resummation [7, 8]. Here we consider the event shape DIS thrust, τ , which is defined in the Breit frame using the momentum of the exchanged γ or Z -boson to determine the z -axis, $q = (0, 0, 0, Q)$. It can be measured solely from events in the current hemisphere where $z > 0$ via $\tau = 1 - (2/Q) \sum_{i \in \mathcal{H}_J} p_{iz}$, thus avoiding the lack of detector coverage in parts of the beam region. The event shape τ also does not suffer from non-global logarithms [8].

Recently an all orders factorization theorem was derived for $d\sigma/d\tau$ [9], which enables higher order perturbative results to be obtained, and a more rigorous treatment of power corrections,

$$\frac{1}{\sigma_0} \frac{d\sigma}{dx dQ^2 d\tau} = Q H(Q, x, \mu) \int dt_B dt_J d^2 p_\perp B_q(t_B, x, \vec{p}_\perp^2, \mu) J_q(t_J - \vec{p}_\perp^2, \mu) S\left(Q\tau - \frac{t_B + t_J}{Q}, \mu\right). \quad (1)$$

Results with a resummation of the singular $\alpha_s^k \ln^j \tau/\tau$ terms at next-to-next-to-leading-log (NNLL) order were given in [9]. Here we extend this analysis to one higher order, N³LL, by exploiting the recent 2-loop calculation of the quark beam function B_q [10, 11], the 2-loop DIS soft function S [12, 13, 14, 15], and known results for the 2-loop hard function H and jet function J_q , and their 3-loop anomalous dimensions [16, 17, 18]. The smaller nonsingular contributions to $d\sigma/d\tau$ are also now known analytically at $\mathcal{O}(\alpha_s)$ [19], while numerical results are available at $\mathcal{O}(\alpha_s^2)$ [20, 21]. Power corrections are encoded by a hadronic matrix element Ω_1 appearing in S , using formalism developed in Refs. [22, 23, 24, 25, 9]. (The DIS thrust τ is equal to the Breit frame 1-jettiness τ_1^b , and hence belongs to the class of 1-jettiness event shapes [26]. Results for other 1-jettiness DIS variables were obtained in Refs. [27, 28, 29, 9], currently up to NNLL order.)

Fits for $\alpha_s(m_Z)$ in the tail region of the DIS τ distribution, should simultaneously fit for the power correction Ω_1 (similar to the highly successful fits for the e^+e^- thrust event shape in [30]). This is facilitated by considering $d\sigma/d\tau$ from multiple x and Q values. Interestingly, the factorization theorem in Eq. (1) remains valid for relatively small x , and the fractional contribution from the nonsingular corrections even decreases with decreasing x , as shown at $\mathcal{O}(\alpha_s)$ in [31].

In Fig. 1 we show the convergence of the DIS thrust cross section and decrease in the perturbative resummation uncertainty when going from NLL to NNLL to N³LL order. Results are displayed for a representative value of x and Q , while cross sections for other values can be found in [31]. In Fig. 2 we show the percent uncertainty of $d\sigma/d\tau$ for various values of x and Q in the region accessible by HERA, demonstrating that the theoretical resummation uncertainties become as low as 2% in accessible regions of the phase space. Values are obtained as the average uncertainty in $d\sigma/d\tau$ in the tail region $0.15 < \tau < 0.35$. In Fig. 3 we show how much the cross section changes with variations of the input parameters $\alpha_s(m_Z)$ and Ω_1 , as well as comparing the $\alpha_s(m_Z)$ sensitivity to the N³LL resummation uncertainties, and to the uncertainties from the NNLO MSTW parton distributions [32]. Figures for other values of x and Q are available in [31]. The degeneracy between $\alpha_s(m_Z)$ and Ω_1 is broken by measurements at multiple Q . The theoretical precision of our N³LL cross section indicates that measurements with 1-2% uncertainty in $\alpha_s(m_Z)$ should now be possible. A measurement of Ω_1 from DIS is also of broader use, since this same Ω_1 parameter occurs in $pp \rightarrow Z + 1\text{-jet}$, where it yields the power correction for the jet-mass that is linear in the jet radius [33].

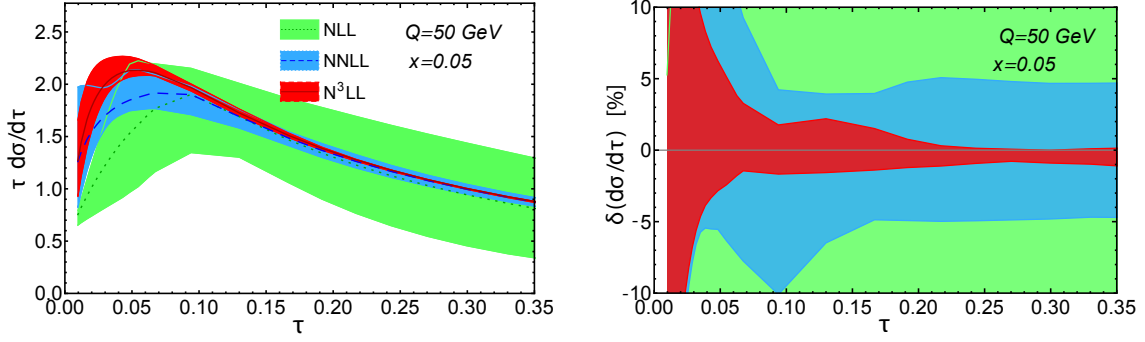


Figure 1: Convergence of the DIS thrust distribution. Results at three orders are shown along with their perturbative uncertainty. Left panel shows $\tau d\sigma/d\tau$. Right panel shows the relative uncertainty for $d\sigma/d\tau$.

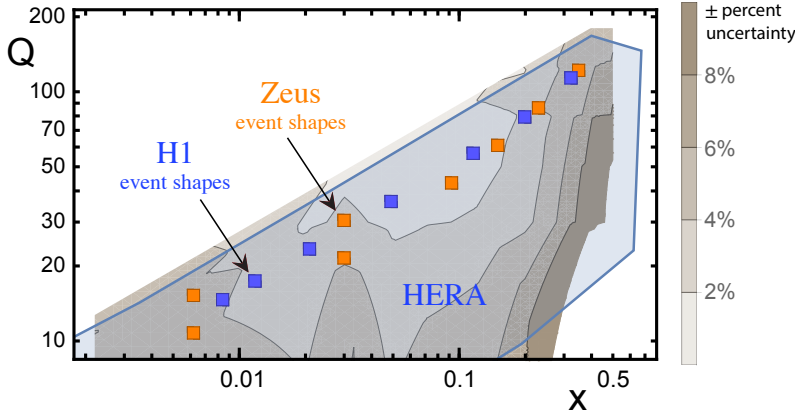


Figure 2: Percent uncertainty of our N^3LL cross section for the region in x and Q accessible at HERA. Uncertainties are for the tail of the DIS thrust distribution which can be used to measure $\alpha_s(m_Z)$. Also shown are the points used for past DIS event shape measurements.

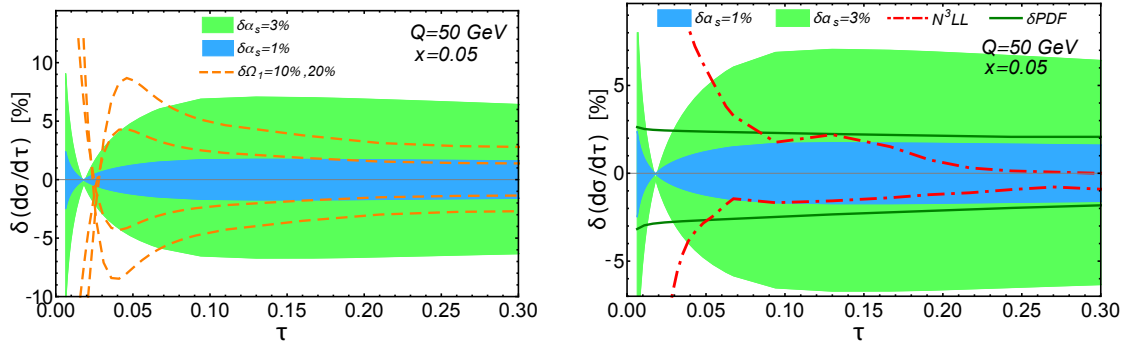


Figure 3: Sensitivity of the DIS thrust cross section to changes in $\alpha_s(m_Z)$ and Ω_1 (left panel) and compared with PDF and N^3LL uncertainties (right panel).

Acknowledgments

The work of IS is supported by the Office of Nuclear Physics of the U.S. Department of Energy

under Contract DE-FG02-94ER40818, and the work of CL and DK by DOE Contract DE-AC52-06NA25396 and by the LDRD office at LANL.

References

- [1] C. Adloff *et al.* [H1 Collaboration], Phys. Lett. B **406**, 256 (1997) [hep-ex/9706002].
- [2] C. Adloff *et al.* [H1 Collaboration], Eur. Phys. J. C **14**, 255 (2000) [Erratum-ibid. C **18**, 417 (2000)] [hep-ex/9912052].
- [3] A. Aktas *et al.* [H1 Collaboration], Eur. Phys. J. C **46**, 343 (2006) [hep-ex/0512014].
- [4] J. Breitweg *et al.* [ZEUS Collaboration], Phys. Lett. B **421**, 368 (1998) [hep-ex/9710027].
- [5] S. Chekanov *et al.* [ZEUS Collaboration], Eur. Phys. J. C **27**, 531 (2003) [hep-ex/0211040].
- [6] S. Chekanov *et al.* [ZEUS Collaboration], Nucl. Phys. B **767**, 1 (2007) [hep-ex/0604032].
- [7] V. Antonelli, M. Dasgupta and G. P. Salam, JHEP **0002**, 001 (2000) [hep-ph/9912488].
- [8] M. Dasgupta and G. P. Salam, JHEP **0208**, 032 (2002) [hep-ph/0208073].
- [9] D. Kang, C. Lee and I. W. Stewart, Phys. Rev. D **88** (2013) 054004 [arXiv:1303.6952 [hep-ph]].
- [10] J. R. Gaunt, M. Stahlhofen and F. J. Tackmann, JHEP **1404** (2014) 113 [arXiv:1401.5478 [hep-ph]].
- [11] J. R. Gaunt and M. Stahlhofen, JHEP **1412** (2014) 146 [arXiv:1409.8281 [hep-ph]].
- [12] D. Kang, O. Z. Labun and C. Lee, Phys. Lett. B **748** (2015) 45 [arXiv:1504.04006 [hep-ph]].
- [13] R. Boughezal, X. Liu and F. Petriello, Phys. Rev. D **91** (2015) 9, 094035 [arXiv:1504.02540 [hep-ph]].
- [14] R. Kelley, M. D. Schwartz, R. M. Schabinger and H. X. Zhu, Phys. Rev. D **84** (2011) 045022 [arXiv:1105.3676 [hep-ph]].
- [15] P. F. Monni, T. Gehrmann and G. Luisoni, JHEP **1108** (2011) 010 [arXiv:1105.4560 [hep-ph]].
- [16] T. Matsuura, S. C. van der Marck and W. L. van Neerven, Nucl. Phys. B **319** (1989) 570.
- [17] T. Becher and M. Neubert, Phys. Lett. B **637** (2006) 251 [hep-ph/0603140].
- [18] S. Moch, J. A. M. Vermaseren and A. Vogt, Nucl. Phys. B **688** (2004) 101 [hep-ph/0403192].
- [19] D. Kang, C. Lee and I. W. Stewart, JHEP **1411** (2014) 132 [arXiv:1407.6706 [hep-ph]].
- [20] S. Catani and M. H. Seymour, Nucl. Phys. B **485** (1997) 291 [Nucl. Phys. B **510** (1998) 503] [hep-ph/9605323].
- [21] D. Graudenz, hep-ph/9710244.
- [22] C. Lee and G. F. Sterman, eConf C **0601121**, A001 (2006) [hep-ph/0603066].
- [23] C. Lee and G. F. Sterman, Phys. Rev. D **75**, 014022 (2007) [hep-ph/0611061].
- [24] A. H. Hoang and I. W. Stewart, Phys. Lett. B **660** (2008) 483 [arXiv:0709.3519 [hep-ph]].
- [25] V. Mateu, I. W. Stewart and J. Thaler, Phys. Rev. D **87**, 014025 (2013) [arXiv:1209.3781 [hep-ph]].
- [26] I. W. Stewart, F. J. Tackmann and W. J. Waalewijn, Phys. Rev. Lett. **105**, 092002 (2010) [arXiv:1004.2489 [hep-ph]].
- [27] Z. -B. Kang, S. Mantry and J. -W. Qiu, Phys. Rev. D **86**, 114011 (2012) [arXiv:1204.5469 [hep-ph]].

- [28] Z. B. Kang, X. Liu and S. Mantry, Phys. Rev. D **90** (2014) 1, 014041 [arXiv:1312.0301 [hep-ph]].
- [29] Z. -B. Kang, X. Liu, S. Mantry and J. -W. Qiu, arXiv:1303.3063 [hep-ph].
- [30] R. Abbate, M. Fickinger, A. H. Hoang, V. Mateu and I. W. Stewart, Phys. Rev. D **83**, 074021 (2011) [arXiv:1006.3080 [hep-ph]].
- [31] D. Kang, C. Lee and I. W. Stewart, "Precision Jet Physics in Deep Inelastic Scattering" for more information see the talk at the Nov. 2014 HERA workshop, "<https://indico.desy.de/contributionDisplay.py?sessionId=10&contribId=33&confId=10523>" and "DIS Event Shape at N3LL" at the DIS 2015 conference, "<https://indico.cern.ch/event/341292/session/14/contribution/24/attachments/670208/921244/DIS2015-kang.pdf>"
- [32] A. D. Martin, W. J. Stirling, R. S. Thorne and G. Watt, Eur. Phys. J. C **63** (2009) 189 [arXiv:0901.0002 [hep-ph]].
- [33] I. W. Stewart, F. J. Tackmann and W. J. Waalewijn, Phys. Rev. Lett. **114** (2015) 9, 092001 [arXiv:1405.6722 [hep-ph]].