

Exclusive production of two and four pions in ultraperipheral, ultrarelativistic collisions

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The cross section for $\pi^+\pi^-$ and $\pi^0\pi^0$ meson pairs production in peripheral ultrarelativistic heavy-ion collisions is calculated in the impact parameter space equivalent photon approximation at the energy available at the CERN Large Hadron Collider, $\sqrt{s_{NN}} = 3.5$ TeV. For the first time the world data for $\gamma\gamma \rightarrow \pi\pi$ are described, both for the total cross section and for the angular distributions simultaneously for $\gamma\gamma \rightarrow \pi^+\pi^-$ and $\gamma\gamma \rightarrow \pi^0\pi^0$ at all experimentally available energies.

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1. Introduction

Ultrarelativistic colliding heavy ions are a source of high-energy $\gamma\gamma$ collisions. Recently we have studied several processes initiated by the photon-photon collisions such as $\rho^0\rho^0$ [1], $\mu^+\mu^-$ [2], $Q\bar{Q}$ [3], $J/\Psi J/\Psi$ [4] and $\pi\pi$ [5, 6]. We have shown there that inclusion of realistic charge form factors, being Fourier transforms of realistic charge distributions, is crucial to estimate reliable nuclear cross sections. Here we consider $PbPb \rightarrow PbPb\pi^+\pi^-$ and $PbPb \rightarrow PbPb\pi^0\pi^0$ reactions. To calculate nuclear cross section, we have to take into account the correct form of the elementary cross section for $\gamma\gamma \rightarrow \pi\pi$ reactions at low, intermediate and high sub-energies.

2. Elementary cross section

The $\gamma\gamma \rightarrow \pi\pi$ reaction is rather complicated. We have tried to understand both $\gamma\gamma \rightarrow \pi^+\pi^-$ and $\gamma\gamma \rightarrow \pi^0\pi^0$ processes simultaneously, starting from the two-pion threshold ($W = 2m_\pi$) up to the maximal experimentally available energy $W_{\gamma\gamma} \approx 6$ GeV. In Fig. 1 we show our model results (left and middle panel) against world data for $\gamma\gamma \rightarrow \pi\pi$. We get a good agreement with all available data for the first time in such a large range of energies. In the right side of Fig. 1 we present a ratio of the neutral and charged pion pair cross sections as a function of $\gamma\gamma$ energy. The upper line represents the handbag model result, which is independent of $z = \cos\theta$ and of the collision energy. Our result, which includes the Brodsky-Lepage pQCD [7] and the handbag mechanism proposed by Diehl, Kroll and Vogt [8]), describes the experimental data measured by the Belle Collaboration.

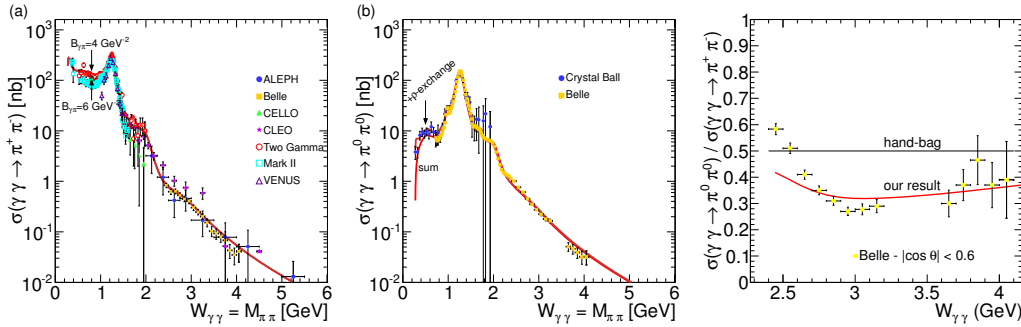


Figure 1: Results of our fit: $\gamma\gamma \rightarrow \pi^+\pi^-$ ($|\cos\theta| < 0.6$) (left panel) and $\gamma\gamma \rightarrow \pi^0\pi^0$ ($|\cos\theta| < 0.8$) (middle panel). The right panel: the ratio of the $\gamma\gamma \rightarrow \pi^0\pi^0$ and $\gamma\gamma \rightarrow \pi^+\pi^-$ cross sections as a function of $\sqrt{s_{\gamma\gamma}}$.

3. Nuclear cross section

The nuclear cross section has been calculated with the help of b -space equivalent photon approximation (EPA). This approach allows to separate ultraperipheral collisions of nuclei ($b > R_1 + R_2 \approx 14$ fm). A compact formula for calculating the total cross section takes the form:

$$\sigma(AA \rightarrow AA\pi\pi; s_{AA}) = \int \hat{\sigma}(\gamma\gamma \rightarrow \pi\pi; W_{\gamma\gamma}) S_{abs}^2(\mathbf{b}) N(\omega_1, \mathbf{b}_1) N(\omega_2, \mathbf{b}_2) \frac{W_{\gamma\gamma}}{2} d^2\mathbf{b}_1 d^2\mathbf{b}_2 dW_{\gamma\gamma} dY_{\pi\pi}.$$

The details of its derivation can be found in our papers [2, 6].

Fig. 2 shows two-dimensional distributions in pseudorapidity of charged (left panel) or neutral pion (right panel) and transverse momentum of one of the pions. With larger $p_{t,\pi}$ values, the pseudorapidity distribution becomes somewhat narrower.

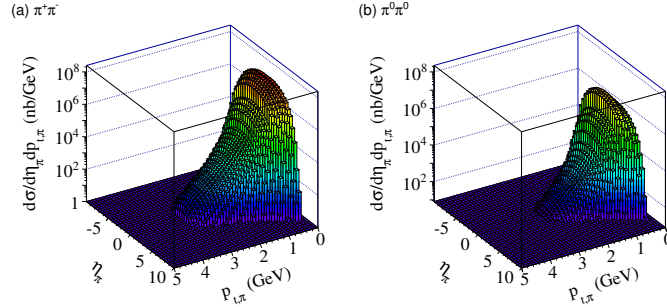


Figure 2: $\frac{d\sigma}{d\eta_{\pi} d p_{t,\pi}}$ for the $^{208}\text{Pb}^{208}\text{Pb} \rightarrow ^{208}\text{Pb}^{208}\text{Pb} \pi^+ \pi^-$ reaction (left panel) and for the $^{208}\text{Pb}^{208}\text{Pb} \rightarrow ^{208}\text{Pb}^{208}\text{Pb} \pi^0 \pi^0$ reaction (right panel) at the LHC energy.

4. Conclusions

We have shown that in order to describe the word data for $\gamma\gamma \rightarrow \pi^+ \pi^-$ and $\gamma\gamma \rightarrow \pi^0 \pi^0$ reactions, one should include several different mechanisms: soft two-pion continuum, several resonances, pion-pion rescattering, pQCD mechanisms proposed by Brodsky and Lepage, as well as the hand-bag mechanism proposed by Diehl, Kroll and Vogt. The energy-dependent cross sections for these two sub-processes have been used next in EPA in the impact parameter space to calculate for the first time corresponding production rate in ultraperipheral ultrarelativistic heavy-ion reactions. In this calculation we have taken into account realistic charge distributions in colliding nuclei.

The $\gamma\gamma \rightarrow \pi^+ \pi^-$ sub-process constitutes a background to the exclusive $AA \rightarrow A\rho^0(\rightarrow \pi^+ \pi^-)A$ process, initiated by photon-pomeron or pomeron-photon sub-processes. We have found that only a part of the dipion invariant mass spectrum associated with $\gamma\gamma$ -collisions can be potentially visible as the cross section for the $AA \rightarrow A\rho^0 A$ reaction is very large [9].

References

- [1] M. Kłusek, A. Szczurek and W. Schäfer, Phys. Lett. **B674** 92 (2009).
- [2] M. Kłusek-Gawenda and A. Szczurek, Phys. Rev. **C82** 014904 (2010).
- [3] M. Kłusek-Gawenda, A. Szczurek, M.V.T. Machado and V.G. Serbo, Phys. Rev. **C83** 024903 (2011).
- [4] S. Baranov, A. Cisek, M. Kłusek-Gawenda, W. Schäfer and A. Szczurek, Eur. Phys. J. **C73** 2335 (2013).
- [5] M. Kłusek-Gawenda and A. Szczurek, Phys. Lett. **B700** 322 (2011).
- [6] M. Kłusek-Gawenda and A. Szczurek, Phys. Rev. **C87** 054908 (2013).
- [7] S.J. Brodsky and G.P. Lepage, Phys. Rev. **D24** 1808 (1981).
- [8] M. Diehl, P. Kroll and C. Vogt, Phys. Lett. **B532** 99 (2002); M. Diehl and P. Kroll, Phys. Lett. **B683** 165 (2010).
- [9] M. Kłusek-Gawenda and A. Szczurek, arXiv: 1309.2463 [nucl-th].