

Light neutralino dark matter in MSSM

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Three dark matter direct detection experiments, DAMA, COGENT and CRESST, have reported a possible signal of WIMP interaction corresponding to very light particles, close to the edge of the XENON 100 and CDMS sensitivity. Imposing the latest constraints from colliders, flavour physics, electroweak precision tests and direct and indirect dark matter searches, we show that viable MSSM scenarios with a light neutralino, in agreement with all the present data, are feasible. An analysis of the characteristics of the resulting scenarios will be presented.

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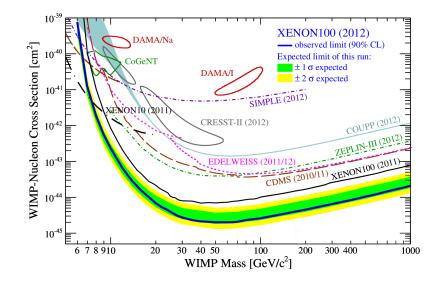


Figure 1: Latest XENON 100 result on spin-independent WIMP-proton scattering cross section [3]. The green/yellow bands correspond to the $1\sigma/2\sigma$ expected sensitivity of the 2012 run. The resulting exclusion limit at 90% C.L. is shown in blue. The exclusion limits from SIMPLE, COUPP, ZEPLIN-III, EDELWEISS, CDMS and the 2011 XENON 100 are displayed for comparison [1]. The regions favoured by DAMA, CoGeNT and CRESST are also shown [4, 5, 6].

1. Introduction

The sensitivity of the dark matter direct detection experiments has strongly improved over the past years [1, 2, 3]. In particular the 2012 XENON 100 result [3] places a 90% C.L. upper limit on the spin-independent WIMP cross section with proton of around 10^{-9} pb for $M_{\text{WIMP}} \simeq$ 100 GeV. On the other hand, three experiments, namely DAMA [4] at the Laboratory Nazionali del Gran Sasso, CoGeNT [5] in the Soudan mine and CRESST [6] also at Gran Sasso, have reported a possible signal of WIMP interaction with matter corresponding to very light particles of masses 5 $< M_{\text{WIMP}} < 15$ GeV, close to the edge of the XENON 100 and CDMS sensitivity. Fig. 1 summarises the sensitivity of the different dark matter experiments.

There is still a large debate on the interpretation of the data and the compatibility of the DAMA, CoGeNT and CRESST results with the exclusion bounds established by the other experiments, however an eventual agreement could be possible only in the low WIMP mass region. In [7] we studied in detail the viability of light neutralino scenarios compatible with all the available data. The main results will be described in the following.

2. pMSSM scans and constraints

We perform flat scans in the phenomenological MSSM (pMSSM) [8] following the procedure described in [9, 10]. The pMSSM with 19 free parameters is the most general CP and R parity conserving MSSM scenario assuming minimal flavour violation and the absence of FCNCs at tree level. The parameters are scanned over in the ranges given in Table 1. In particular, we compute the SUSY spectra with SOFTSUSY [11], flavour physics observables and relic density with SuperIso Relic [12], and the dark matter scattering cross sections with Micromegas [13].

Parameter	Range	Parameter	Range
$\tan\beta$	[1, 60]	$M_{\tilde{e}_L} = M_{\tilde{\mu}_L}$	[0, 2500]
M _A	[50, 2000]	$M_{ ilde{e}_R} = M_{ ilde{\mu}_R}$	[0, 2500]
M_1	[-300, 300]	$M_{ ilde{ au}_L}$	[0, 2500]
<i>M</i> ₂	[-650, 650]	$M_{ ilde{ au}_R}$	[0, 2500]
<i>M</i> ₃	[0, 2500]	$M_{\tilde{q}_{1L}} = M_{\tilde{q}_{2L}}$	[0, 2500]
$A_d = A_s = A_b$	[-10000, 10000]	$M_{ ilde{q}_{3L}}$	[0, 2500]
$A_u = A_c = A_t$	[-10000, 10000]	$M_{\tilde{u}_R} = M_{\tilde{c}_R}$	[0, 2500]
$A_e = A_\mu = A_\tau$	[-10000, 10000]	$M_{\tilde{t}_R}$	[0, 2500]
μ	[-3000, 3000]	$M_{\tilde{d}_R} = M_{\tilde{s}_R}$	[0, 2500]
		$M_{\tilde{b}_R}$	[0, 2500]

Table 1: pMSSM parameter ranges used in the scans (in GeV when applicable).

$2.16 \times 10^{-4} < \mathrm{BR}(B \to X_s \gamma) < 4.93 \times 10^{-4}$		
$BR(B_s \to \mu^+ \mu^-) < 5.0 \times 10^{-9}$		
$0.56 < \mathbf{R}(B \to \tau \nu) < 2.70$		
$4.7 \times 10^{-2} < \mathrm{BR}(D_s \to \tau \nu) < 6.1 \times 10^{-2}$		
$2.9 \times 10^{-3} < \mathrm{BR}(B \to D^0 \tau \nu) < 14.2 \times 10^{-3}$		
$0.985 < R_{\mu 23}(K \to \mu \nu) < 1.013$		
$-2.4 imes 10^{-9} < \delta a_{\mu} < 4.5 imes 10^{-9}$		

Table 2: Low energy constraints applied in our analysis.

We focus on light neutralino masses in the range $5 < M_{\tilde{\chi}_1^0} < 40$ GeV and spin-independent scattering cross sections with proton in the region where the dark matter direct detection data can be reconciled, $10^{-7} < \sigma_{\chi p}^{SI} < 10^{-3}$ pb. For the relic density, based on the WMAP data [14] we consider a tight interval at 95% C.L. including theoretical uncertainties:

$$0.068 < \Omega_{\gamma} h^2 < 0.155 , \qquad (2.1)$$

as well as a loose interval:

$$10^{-4} < \Omega_{\chi} h^2 < 0.155 , \qquad (2.2)$$

which takes into account the possibility of multiple species dark matter or modified cosmological models [15, 16]. We also apply the constraints from flavour physics observables and muon anomalous magnetic moment given in Table 2, as well as recent constraints from Higgs, SUSY and monojet searches at LEP, Tevatron and LHC. More details can be found in [7].

3. Light neutralinos in the pMSSM

Using a large statistics of more than one billion points in general scans, and more than one billion extra points in specific scans, we were able to identify three classes of pMSSM solutions

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satisfying all the constraints described above, where

- class i) the NLSP is a slepton slightly above the LEP limit, with a neutralino of about 30 GeV - class ii) the lightest chargino is degenerate with the $\tilde{\chi}_1^0$, often with a compressed gaugino spectrum and light Higgs bosons

- class iii) a scalar quark is degenerate with the $\tilde{\chi}_1^0$ while the other supersymmetric particles are relatively heavy. Imposing a light CP-even Higgs mass of ~ 126 GeV, the only possible squarks for class iii) are those of the first and second generations or the lightest scalar bottom quark \tilde{b}_1 .

Scenario i) is only of limited interest, since it leads to large scattering cross sections pointing to a region well inside the exclusion contours by the CDMS and XENON experiments where reconciliation between the different direct detection experiment results seems difficult.

In Scenario ii), the spin independent χ -*p* scattering cross section is relatively small (~ 10⁻⁶ – 10⁻⁷ pb), and because the production cross section of $\chi_1^+\chi_1^-$ and $\chi_2^0\chi_1^0$ in this scenario at LEP-2 is large, the combination of the LEP-2 experiment results [17] could disfavour this scenario.

Scenario iii) corresponds to large spin independent χ -*p* scattering cross sections associated to small neutralino masses of about 10 GeV, and is therefore the most interesting scenario. The couplings of a light squark to the *Z* and *h* bosons are in general large, and the *Z* decay width into squarks would exclude this scenario unless the squark decouples from the *Z*. This is the case for specific values of the squark mixing angle. For this reason, because the first and second generation squarks do not mix, they are excluded by LEP-1 data, and the only viable scenarios correspond to a scalar bottom \tilde{b}_1 nearly degenerate with the lightest neutralino. If the right-handed bottom squark is very light, the mixing angle θ_b is large and the squark decouples from the *Z*. In addition, higher order SUSY corrections further decrease the \tilde{b}_1 mass for specific values of the other parameters. In this context, it possible to find pMSSM solutions with light neutralinos which are compatible with the LEP-1 constraints. However, the decay rate $h \rightarrow \tilde{b}_1 \tilde{b}_1$ can become important in such scenarios, resulting in the reduction of the other Higgs decay rates, which is disfavoured by the recent Higgs data [18, 19]. Nevertheless, there exist model points for which the branching ratio of $h \rightarrow \tilde{b}_1 \tilde{b}_1$ and the *Z* decay width to $\tilde{b}_1 \tilde{b}_1$ are simultaneously small, making this scenario compatible with the LHC Higgs search data.

To assess the viability of Scenario iii) in view of the dark matter indirect detection data, we consider the latest constraints by Fermi-LAT [20]. We show that the points with degenerate \tilde{b}_1 have neutralino annihilation cross sections times relative velocity of about one order of magnitude below the current Fermi-LAT limits, so that Scenario iii) is also compatible with the dark matter indirect detection limits.

In Fig. 2, we present the distribution of the points passing the tight relic density bounds (2.1), the loose relic density bounds (2.2), and without relic density constraint. A comparison of the three plots shows that the relic density lower bound decreases the overall statistics, but also removes scenarios with a degenerate sbottom and neutralino masses above 20 GeV. This is due to the fact that points with a very small relic density generally have a small splitting. However, in order to retrieve a relic density in the interval (2.1), the splitting of the NLSP with the neutralino should not be too small. In addition, large direct detection scattering cross sections disfavour large splittings. As a consequence, combining the relic density with direct search limits points to only a small viable region where all the constraints can be fulfilled simultaneously. We also note that relaxing the relic density constraint would restore the viability of the three classes of scenarios.

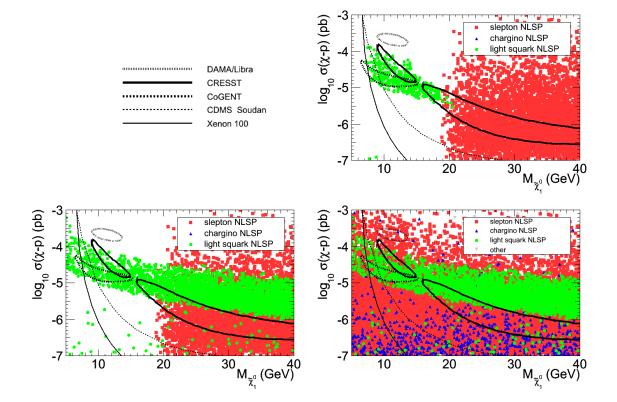


Figure 2: Spin independent χ -p scattering cross section as a function of the $\tilde{\chi}_1^0$ mass. The points plotted here pass all the constraints, including the tight relic density constraint (upper right), the loose relic density constraint (lower left) and no relic density constraint (lower right). The regions favoured by CRESST, CoGeNT and DAMA, as well as the exclusion limit by CDMS and XENON 100 are also shown for comparison. The red squares correspond to scenarios with a slepton NLSP with a mass slightly above the LEP limits (class i), the blue triangles to cases with a chargino NLSP (class ii), and the green points to the class where a sbottom is nearly degenerate with the lightest neutralino (class iii).

4. Conclusion

We have investigated solutions in the generic minimal supersymmetric extension of the Standard Model with light neutralinos, considering the constraints from flavour physics, low energy, LEP, Tevatron and LHC data. Using very high statistics flat scans of the pMSSM we found three different scenarios with simultaneously very light neutralinos and large direct detection scattering cross sections. Imposing the LEP and LEP-2 limits, and at the same time the relic density constraint, only one viable scenario with a very light scalar bottom quark \tilde{b}_1 nearly degenerate with the neutralino remains. We checked that such scenario is compatible with the monojet searches at the LHC and direct detection data from DAMA, CoGeNT and CRESST in a region which could be reconciled with the CDMS and XENON limits, as well as indirect constraints from Fermi-LAT. Relaxing the relic density constraints would however allow to restore the viability of the two other classes of scenarios with chargino and slepton NLSPs.

References

Z. Ahmed *et al.* [CDMS-II Collaboration], Science **327** (2010) 1619 [arXiv:0912.3592]; Phys. Rev. Lett. **106** (2011) 131302 [arXiv:1011.2482];
 E. Armengaud *et al.* [EDELWEISS Collaboration], Phys. Lett. B **702** (2011) 329 [arXiv:1103.4070];

J. Angle *et al.* [XENON 10 Collaboration], Phys. Rev. Lett. **107** (2011) 051301 [arXiv:1104.3088];
M. Felizardo *et al.* [SIMPLE Collaboration], Phys. Rev. Lett. **108** (2012) 201302 [arXiv:1106.3014];
E. Behnke *et al.* [COUPP Collaboration], Phys. Rev. D **86** (2012) 052001 [arXiv:1204.3094];
D. Y. .Akimov *et al.* [ZEPLIN Collaboration], Phys. Lett. B **709** (2012) 14 [arXiv:1110.4769];
E. Armengaud *et al.* [EDELWEISS Collaboration], Phys. Rev. D **86** (2012) 051701

[arXiv:1207.1815].

- [2] E. Aprile et al. [XENON 100 Collaboration], Phys. Rev. Lett. 107 (2011) 131302 [arXiv:1104.2549].
- [3] E. Aprile et al. [XENON 100 Collaboration], Phys. Rev. Lett. 109 (2012) 181301 [arXiv:1207.5988].
- [4] C. Savage, G. Gelmini, P. Gondolo and K. Freese, JCAP 0904 (2009) 010 [arXiv:0808.3607].
- [5] C. E. Aalseth et al. [CoGeNT Collaboration], Phys. Rev. Lett. 106 (2011) 131301 [arXiv:1002.4703].
- [6] G. Angloher et al. [CRESST-II Collaboration], Eur. Phys. J. C 72 (2012) 1971 [arXiv:1109.0702].
- [7] A. Arbey, M. Battaglia and F. Mahmoudi, Eur. Phys. J. C 72 (2012) 2169 [arXiv:1205.2557].
- [8] A. Djouadi et al. [MSSM Working Group Collaboration], hep-ph/9901246.
- [9] A. Arbey, M. Battaglia and F. Mahmoudi, Eur. Phys. J. C 72 (2012) 1847 [arXiv:1110.3726].
- [10] A. Arbey, M. Battaglia and F. Mahmoudi, Eur. Phys. J. C 72 (2012) 1906 [arXiv:1112.3032].
- [11] B. C. Allanach, Comput. Phys. Commun. 143 (2002) 305 [hep-ph/0104145].
- [12] F. Mahmoudi, Comput. Phys. Commun. 178 (2008) 745 [arXiv:0710.2067]; Comput. Phys. Commun. 180 (2009) 1579 [arXiv:0808.3144]; Comput. Phys. Commun. 180 (2009) 1718;
 A. Arbey and F. Mahmoudi, Comput. Phys. Commun. 181 (2010) 1277 [arXiv:0906.0369]; Comput. Phys. Commun. 182 (2011) 1582.
- [13] G. Bélanger *et al.*, Comput. Phys. Commun. **176** (2007) 367 [hep-ph/0607059]; Comput. Phys. Commun. **180** (2009) 747 [arXiv:0803.2360].
- [14] E. Komatsu et al. [WMAP Collaboration], Astrophys. J. Suppl. 192 (2011) 18 [arXiv:1001.4538].
- [15] A. Arbey and F. Mahmoudi, Phys. Lett. B 669 (2008) 46 [arXiv:0803.0741].
- [16] A. Arbey and F. Mahmoudi, JHEP 1005 (2010) 051 [arXiv:0906.0368].
- [17] LEP2 SUSY Working Group, Note LEPSUSYWG/02-04.1.
- [18] A. Arbey, M. Battaglia, A. Djouadi, F. Mahmoudi and J. Quevillon, Phys. Lett. B 708 (2012) 162 [arXiv:1112.3028].
- [19] A. Arbey, M. Battaglia, A. Djouadi and F. Mahmoudi, JHEP 1209 (2012) 107 [arXiv:1207.1348].
- [20] M. Ackermann *et al.* [Fermi-LAT Collaboration], Phys. Rev. Lett. **107** (2011) 241302 [arXiv:1108.3546].