

## Observation of a Narrow Near-Threshold Structure in the $J/\psi\phi$ Mass Spectrum in $B^+ \rightarrow J/\psi\phi K^+$ Decays

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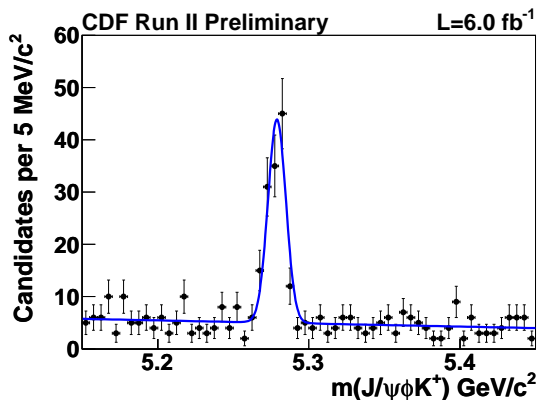
Observation is reported for a structure near the  $J/\psi\phi$  threshold in  $B^+ \rightarrow J/\psi\phi K^+$  decays produced in  $\bar{p}p$  collisions at  $\sqrt{s} = 1.96$  TeV with a statistical significance of beyond 5 standard deviations. There are  $19 \pm 6$  events observed for this structure at a mass of  $4143.4_{-3.0}^{+2.9}(\text{stat}) \pm 0.6(\text{syst})$  MeV/ $c^2$  and a width of  $15.3_{-6.1}^{+10.4}(\text{stat}) \pm 2.5(\text{syst})$  MeV/ $c^2$ , which are consistent with the previous measurements reported as evidence of the  $Y(4140)$ .

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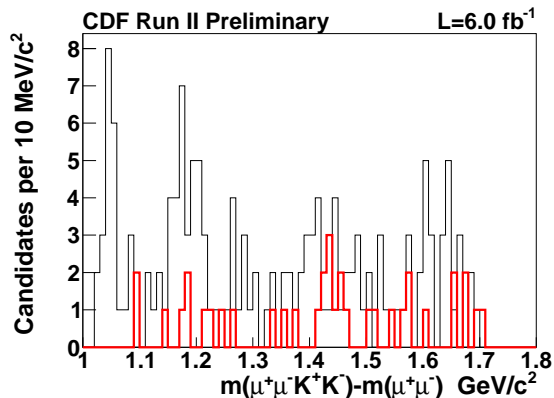
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**Figure 1:** The mass distribution of  $J/\psi\phi K^+$ ; the solid line is a fit to the data with a Gaussian signal function and linear background function.

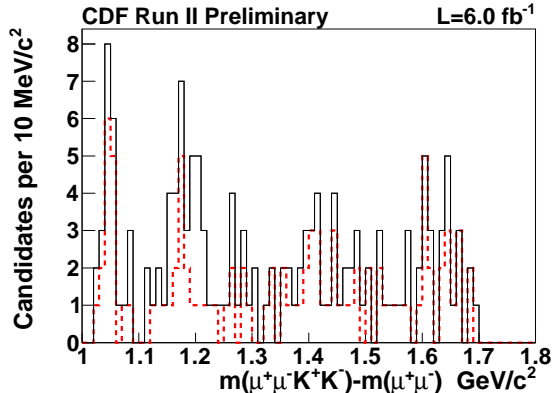


**Figure 2:** The mass difference,  $\Delta M$ , between  $\mu^+\mu^-K^+K^-$  and  $\mu^+\mu^-$ , in the  $B^+$  mass window is shown as the black histogram. The red histogram is the  $\Delta M$  distribution from the data in the B sideband.

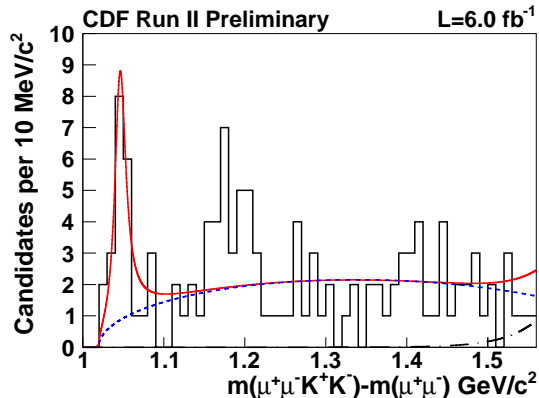
Recently, evidence has been reported by CDF for a narrow structure near the  $J/\psi\phi$  threshold, named  $Y(4140)$ , in  $B^+ \rightarrow J/\psi\phi K^+$  decays produced in  $\bar{p}p$  collisions at  $\sqrt{s} = 1.96$  TeV [1]. The structure is the first charmonium-like structure decaying into two heavy quarkonium states ( $c\bar{c}$  and  $s\bar{s}$ ) which is a candidate for exotic mesons [2]. In this note, we report an update on a search for structures in the  $J/\psi\phi$  system produced in exclusive  $B^+ \rightarrow J/\psi\phi K^+$  decays with  $J/\psi \rightarrow \mu^+\mu^-$  and  $\phi \rightarrow K^+K^-$ . This analysis is based on a sample of  $\bar{p}p$  collision data at  $\sqrt{s} = 1.96$  TeV with an integrated luminosity of about  $6.0 \text{ fb}^{-1}$  collected by the CDF II detector at the Tevatron. The CDF II detector has been described in detail elsewhere [3]. In this analysis,  $J/\psi \rightarrow \mu^+\mu^-$  events are recorded using a dedicated three-level dimuon trigger.

The invariant mass of  $J/\psi\phi K^+$  in the current dataset, which includes those used in the previous analysis after applying the same requirements used in the previous analysis [1], is shown in Fig. 1. A fit with a Gaussian signal function with its rms fixed to the value  $5.9 \text{ MeV}/c^2$  obtained from Monte Carlo (MC) simulation [4] and a linear background function to the mass spectrum of  $J/\psi\phi K^+$  returns a  $B^+$  signal of  $115 \pm 12(\text{stat})$  events. For a luminosity increase by a factor of 2.2, the yield increase was 1.53, reduced by trigger rate-limitation at higher average luminosity. We then select  $B^+$  signal candidates with a mass within  $3\sigma$  ( $17.7 \text{ MeV}/c^2$ ) of the nominal  $B^+$  mass. We define those events with a mass within  $[-9, -6]\sigma$  or  $[6, 9]\sigma$  of nominal  $B$  mass as B sideband events. Fig. 2 shows the mass difference,  $\Delta M = m(\mu^+\mu^-K^+K^-) - m(\mu^+\mu^-)$ , for events in the  $B^+$  mass window as well as in the B sideband in our data sample. The comparison of the  $\Delta M$  distribution in the B mass window for the dataset used in this analysis ( $6.0 \text{ fb}^{-1}$ ) and the dataset used in the previous analysis ( $2.7 \text{ fb}^{-1}$  [1]) is shown in Figure 3.

The same model is used to examine the  $Y(4140)$  structure as described in reference [1]. We model the enhancement by an  $S$ -wave relativistic BW function [5] convoluted with a Gaussian resolution function with the r.m.s. fixed to  $1.7 \text{ MeV}/c^2$  obtained from MC, and use three-body phase space [6] to describe the background shape. Even though we exclude the high mass region to avoid the  $B_s$  contamination, there is still a small contribution in the region of interest. We obtained the  $\Delta M$  shape from  $B_s$  contamination and fix the  $\Delta M$  shape obtained from  $B_s$  MC simulation, and



**Figure 3:** The  $\Delta M$  distribution in the  $B$  mass window for the data used in the current analysis ( $6.0 \text{ fb}^{-1}$ ) is shown as the black histogram, and the same distribution for the data in the previous analysis ( $2.7 \text{ fb}^{-1}$  [1]) is shown as the red dashed histogram.



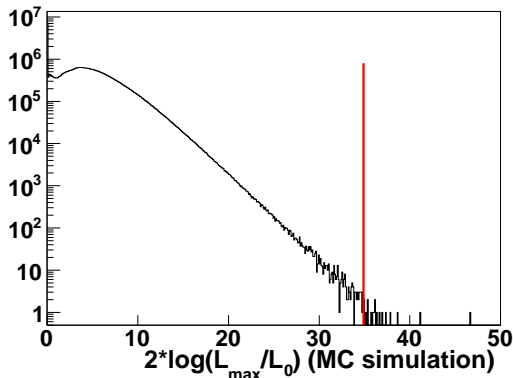
**Figure 4:** The mass difference,  $\Delta M$ , between  $\mu^+\mu^-K^+K^-$  and  $\mu^+\mu^-$ , in the  $B^+$  mass window is shown as a solid black histogram for the data. The dotted curve is the predicted three-body phase space background contribution, the dash-dotted curve is the predicted  $B_s$  contamination (fixed to 3.3), and the solid red curve is the total unbinned fit where the signal PDF is an S-wave Breit-Wigner convoluted with the resolution ( $1.7 \text{ MeV}/c^2$ ).

the yield to 3.3, scaled from the  $B_s \rightarrow J/\psi\phi$  yield in data. An unbinned likelihood fit to the  $\Delta M$  distribution, as shown in Fig. 4, returns a yield of  $19 \pm 6$  events, a  $\Delta M$  of  $1046.7^{+2.9}_{-3.0} \text{ MeV}/c^2$ , and a width of  $15.3^{+10.4}_{-6.1} \text{ MeV}/c^2$ .

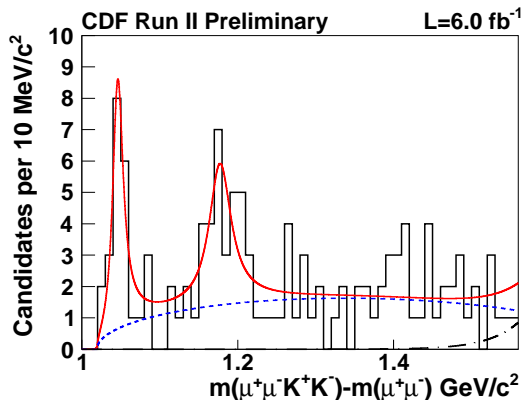
We use the same simulation process as in Reference [1], based on a pure three-body phase space background shape to determine the significance of the  $Y(4140)$  structure. We performed a total of 84 million simulations and found 19 trials with a  $\sqrt{-2\ln(\mathcal{L}_0/\mathcal{L}_{max})}$  value greater than or equal to the value obtained in the data (5.9), as shown in Fig. 5, where  $\mathcal{L}_0$  and  $\mathcal{L}_{max}$  are the likelihood values for the null hypothesis fit and signal hypothesis fit. The resulting  $p$ -value is  $2.3 \times 10^{-7}$ , corresponding to a significance of greater than  $5.0\sigma$ .

The mass of this structure, including systematic uncertainty, is measured as  $4143.4^{+2.9}_{-3.0}(\text{stat}) \pm 0.6(\text{syst}) \text{ MeV}/c^2$  after including the world-average  $J/\psi$  mass. The relative efficiency between  $B^+ \rightarrow Y(4140)K^+, Y(4140) \rightarrow J/\psi\phi$  and  $B^+ \rightarrow J/\psi\phi K^+$  is 1.1 assuming  $Y(4140)$  as an S-wave structure and  $B^+$  phase space decays. Thus the relative branching fraction between  $B^+ \rightarrow Y(4140)K^+, Y(4140) \rightarrow J/\psi\phi$  and  $B^+ \rightarrow J/\psi\phi K^+$  including systematics is  $0.149 \pm 0.039(\text{stat}) \pm 0.024(\text{syst})$ .

An further excess above the three-body phase space background shape appears at approximately  $1.18 \text{ GeV}/c^2$  in Fig. 1 (b). Since the significance of  $Y(4140)$  is greater than  $5\sigma$ , we fit to the data assuming two structures at  $\Delta M$  of 1.05 and  $1.18 \text{ GeV}/c^2$  as shown in Fig. 6. The fit to the data with the same requirements as in the previous paper [1] returns a yield of  $20 \pm 5$  events, a  $\Delta M$  of  $1046.7^{+2.8}_{-2.9} \text{ MeV}/c^2$ , and a width of  $15.0^{+8.5}_{-5.6} \text{ MeV}/c^2$  for the  $Y(4140)$ , which are consistent with the values returned from a  $Y(4140)$ -only signal fit. The fit returns a yield of  $22 \pm 8$  events, a  $\Delta M$  of  $1177.7^{+8.4}_{-6.7} \text{ MeV}/c^2$ , and a width of  $32.3^{+21.9}_{-15.3} \text{ MeV}/c^2$  for the structure around  $\Delta M$  of  $1.18 \text{ GeV}/c^2$ . The significance of the second structure, determined by a similar simulation is  $3.1\sigma$ .



**Figure 5:** Distribution of  $-2\ln(\mathcal{L}_0/\mathcal{L}_{max})$  for 84 million simulation trials. The  $p$ -value obtained from counting is  $2.3 \times 10^{-7}$ , corresponding to a significance of  $5.0\sigma$ .



**Figure 6:** The mass difference,  $\Delta M$ , between  $\mu^+\mu^-K^+K^-$  and  $\mu^+\mu^-$ , in the  $B^+$  mass window. The dotted curve is the background contribution, the dash-dotted curve is the  $B_s$  contamination, and the red solid curve is the total unbinned fit assuming two structures.

In summary, the growing  $B^+ \rightarrow J/\psi\phi K^+$  sample at CDF enables us to observe the  $Y(4140)$  structure [1] with a significance greater than  $5\sigma$ . Assuming an  $S$ -wave relativistic BW, the mass and width of this structure, including systematic uncertainties, are measured to be  $4143.4^{+2.9}_{-3.0}(\text{stat}) \pm 0.6(\text{syst}) \text{ MeV}/c^2$  and  $15.3^{+10.4}_{-6.1}(\text{stat}) \pm 2.5(\text{syst}) \text{ MeV}/c^2$ , respectively. The relative branching fraction between  $B^+ \rightarrow Y(4140)K^+$ ,  $Y(4140) \rightarrow J/\psi\phi$  and  $B^+ \rightarrow J/\psi\phi K^+$  including systematics is  $0.149 \pm 0.039(\text{stat}) \pm 0.024(\text{syst})$ . We also find evidence at  $3.1\sigma$  level for a second structure with a mass of  $4274.4^{+8.4}_{-6.7}(\text{stat}) \text{ MeV}/c^2$ , a width of  $32.3^{+21.9}_{-15.3}(\text{stat}) \text{ MeV}/c^2$  and a yield of  $22 \pm 8$ .

## References

- [1] T. Aaltonen *et al.* (CDF Collaboration), Phys. Rev. Lett. **102**, 242002 (2009).
- [2] X. Liu and S. Zhu, Phys. Rev. D **80**, 017502 (2009); N. Mahajan, Phys. Lett. B **679**, 228 (2009); Z. Wang, Eur. Phys. J. C **63**, 115 (2009); T. Branz *et al.*, Phys. Rev. D **80**, 054019 (2009); R. Albuquerque *et al.*, Phys. Lett. B **678**, 186 (2009); X. Liu, Phys. Lett. B **680**, 137 (2009); G. Ding, Eur. Phys. J. C **64**, 297 (2009); J. Zhang and M. Huang, Phys. Rev. D **80**, 056004 (2009); E. Beveren and G. Rupp, arXiv:0906.2278 [hep-ph]; F. Stancu, J. Phys. G **37**, 075017 (2010); T. Branz *et al.*, arXiv:1001.3959 [hep-ph]; K. Yamada, arXiv:1002.0410 [hep-ph].
- [3] D. Acosta *et al.* (CDF Collaboration), Phys. Rev. D **71**, 032001 (2005); A. Abulencia *et al.* (CDF Collaboration), Phys. Rev. Lett. **97**, 242003 (2006).
- [4] A. Abulencia *et al.* (CDF Collaboration), Phys. Rev. Lett. **96**, 082002 (2006); T. Aaltonen *et al.* (CDF Collaboration), Phys. Rev. Lett. **100** 182002 (2008).
- [5]  $\frac{dN}{dm} \propto \frac{m\Gamma(m)}{(m^2-m_0^2)^2+m_0^2\Gamma^2(m)}$ , where  $\Gamma(m) = \Gamma_0 \frac{q}{q_0} \frac{m_0}{m}$ , and the 0 subscript indicates the value at the peak mass.
- [6] C. Amsler *et al.* (Particle Data Group), Phys. Lett. B **667**, 1 (2008).