

SEARCHES FOR A HIGH MASS STANDARD MODEL HIGGS BOSON AT THE TEVATRON

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Higgs boson searches are commonly considered one of the main objectives of particle physics nowadays. The latest results obtained by the CDF and DØ collaborations when searching for Higgs boson decaying into a W-boson pair, currently the most sensitive channel for masses greater than 130 GeV, are presented here. The presented results are based on an integrated luminosity that ranges from 2.0 to 5.4 fb⁻¹. No significant excess over expected background is observed and the 95% CL limits are set for a Standard Model (SM) Higgs boson for different mass hypotheses ranging from 100 GeV to 200 GeV. The combination of CDF and DØ results is also presented, which excludes a SM Higgs boson in the 162 < m_H < 166 GeV mass range.

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1. Higgs searches at the Tevatron

Hadron colliders collision data can be used to probe Standard Model (SM) theory in several sectors. Data collected at the $p\bar{p}$ Tevatron collider operating at $\sqrt{s} = 1.96$ TeV has been useful for both re-establish or improve previous important measurements and to make new important discoveries, looking for processes whose cross section range by several orders of magnitude. CDF and DØ experiments[1] [2] are collecting data since 2002 and nowadays they have around 6.3 fb^{-1} of good data available for analysis. The Higgs searches presented here are based on an integrated luminosity that goes from 2.0 to 5.4 fb^{-1} depending on the specific analysis.

According to SM, Higgs boson is produced at Tevatron energies by four main production mechanisms, for a total cross-section of about 0.6 pb^1 . The dominant mechanism is via gluon fusion and fermionic loop ($\sim 86\%$), then there is the production with a W or Z boson ($\sim 15\%$) and finally by vector-boson fusion (VBF).

SM Higgs decay depends on its mass m_H . For $m_H > 135$ GeV the main decay channel is a W -boson pair, and searches of this final state at Tevatron are commonly defined as *High Mass Higgs Searches*.

CDF and DØ look for both W s decaying leptonically, selecting events with two electrons or two muons or one electron and one muon for a total branching ratio of the WW pair of $\sim 6\%$, including the leptonic decay of the τ . The larger branching ratio decay in hadrons is not used due to high level of QCD background. The di-lepton final state offer a clean signature at hadron colliders and the trigger cross section is under control also at high instantaneous luminosity.

The main background consists of Drell-Yan (DY) events, which have the same final visible state of the signal (two leptons) but no real missing energy and can be reduced requiring it. Heavy di-boson production (WW , WZ , ZZ) is the most important not reducible background for these searches. $t\bar{t}$ is also an important background, in particular for events containing jets. Finally instrumental backgrounds arise from $W/Z + \gamma$ or *jets* events where the photon or the jet is misidentified as a lepton. Most of these processes are modelled using PYTHIA Monte Carlo and a GEANT-based simulation of the detectors. An important exception is the WW background: CDF models NLO effects using a pure NLO simulation, namely MC@NLO, while DØ uses Sherpa to model the p_T of the WW system[3]. These predictions are then normalized to NNLO cross section calculations for $WH, ZH, t\bar{t}$ processes, NLO for $VBF, WW, WZ, ZZ, W\gamma$. The gluon fusion signal process has been simulated using the most recent calculation available[4] which uses the recent MSTW2008 parton density functions (pdf) set[5]. Data-driven methods are used in order to model the contribution of the instrumental background.

2. CDF and DØ experiments at Tevatron collider

CDF and DØ experiments are multi-purpose cylindrical detectors. The azimuthal angle ϕ and the pseudo-rapidity $\eta \equiv -\ln[\tan(\theta/2)]$, where θ is the polar angle defined relative to the proton beam axis are the basis of the coordinate system. CDF consists of a solenoidal charged particle spectrometer which includes a silicon micro-strip detector array surrounded by a cylindrical drift

¹Higgs boson properties depend on its mass, which is not predicted inside the SM framework; here and whenever not explicitly mentioned in the following, numbers refer as an example to a Higgs boson of mass $m_H = 160$ GeV.

chamber in a 1.4 T axial magnetic field; outside of the tracking chambers, the energies of electrons and jets are measured with segmented sampling calorimeters; the outer-most detectors are layers of steel instrumented with planar drift chambers and scintillators used for muon identification. DØ detector has a central tracking system which consists of a silicon micro-strip tracker and a central fiber tracker, both located within a 2 T axial magnetic field; a liquid-argon/uranium calorimeter surrounds the central tracking system outside of which we find layers of tracking detectors and scintillators inside 1.8 T toroid magnets for muon identification. Both detectors are described in detail elsewhere[1] [2].

The transverse energy, E_T , is defined to be $E \sin \theta$, where E is the energy associated with a calorimeter tower or cluster, and p_T is the component of the track momentum transverse to the beam line. The missing transverse energy vector \vec{E}_T , is defined as the opposite of the vector sum of the E_T of all calorimeter towers, eventually corrected for the p_T of leptons candidates which do not deposit all of their energy in the calorimeter.

3. Analysis description

To select signal events both collaborations require two high- p_T opposite sign isolated leptons. In order to increase acceptance to Higgs events, dedicated analyses also look for final state containing two leptons with the same charge; they will be briefly discussed in section 3.3. CDF requires the first (second) lepton to have p_T greater than 20 (10) GeV, while DØ asks both leptons to have p_T (E_T) greater than 10 (15) GeV for muons (electrons). A significant transverse missing energy \vec{E}_T is then required to reduce DY. The invariant mass of the lepton pair must be greater than 16(15) GeV at CDF(DØ) in order to suppress heavy flavor decays and DY events.

The signal/background ratio is low for a counting experiment to discriminate background-only from signal plus background hypotheses. Other kinematical properties are used to enhance the separation. The strongest discriminant is the opening angle between final state leptons; in signal events leptons come from the decay of spin-1 particles which are from spin-0 Higgs boson, which implies, given the fixed helicity for neutrinos, that leptons will tend to go in the same direction. For the background this spin correlation does not exist and final state leptons tend to be back-to-back. To exploit this and other kinematical signal properties both collaborations use multivariate techniques. An Artificial Neural Network (NN) is trained to separate signal from background for each different Higgs mass hypothesis. Moreover, in order to exploit different signal and background composition, the selected sample is divided in several analysis channels.

3.1 DØ analysis

DØ collaboration separate the analysis depending on the flavor of the final state leptons: ee , $e\mu$, $\mu\mu$, with a collected integrated luminosity of 5.4 fb^{-1} . Input to the NNs can be classified in three different types: lepton specific variables (e.g. p_T of the leptons), kinematic properties of the whole event (e.g. \vec{E}_T) or angular variables (e.g. $\Delta\phi(\text{leptons})$).

No significant excess over predicted background is observed in DØ data and 95% CL upper limits on the production cross section of a SM Higgs boson are set using a modified frequentist approach (CL_s)[9], which are 80% higher than the predicted SM cross section for $m_H = 160 \text{ GeV}$ [7].

3.2 CDF analysis

CDF collaboration separates the 4.8 fb^{-1} selected sample depending on jet multiplicity, optimizing different NNs for each sample. Jets are requested to have $E_T > 15 \text{ GeV}$, $|\eta| < 2.5$. Input variables to the NNs are analogous to $D\emptyset$. Events with no jets have signal contribution only from gluon fusion and the dominant background is WW production. Events with one jet have additional signal contribution from associate Higgs production and VBF and WW still remain the main source of background. Finally, signal events with two or more jets are dominated by WH, ZH and VBF production mechanisms and the dominant background comes from $t\bar{t}$.

No significant excess over predicted background is observed in CDF data and 95% CL upper limit on the production cross section of a SM Higgs boson are set using a Bayesian technique[9], which for $m_H = 160 \text{ GeV}$ are 50% higher than the predicted SM cross section[6].

3.3 Same sign analyses

In order to increase the acceptance to Higgs events, both collaborations also perform searches with two same charge final state leptons. The main signal contribution comes from $WH \rightarrow WWW \rightarrow l^\pm l^\pm + X$, where one of the leptons comes from the W boson produced in association with the Higgs. For these searches the main backgrounds are instrumental backgrounds coming from charge mis-identification or jets faking a lepton signature. CDF uses 4.8 fb^{-1} and an analysis technique similar to the opposite sign analysis to set 95% CL upper limits on the production of a SM Higgs boson that are 5.7 times the expected SM cross section[6] (σ_{SM}^H). $D\emptyset$ has similar analysis which uses 3.6 fb^{-1} of data and set limits up to $18.4 \cdot \sigma_{SM}^H$ [8].

3.4 CDF low M_{ll} analysis

To further increase signal acceptance CDF considers separately events which fail only the $M_{ll} > 16 \text{ GeV}/c^2$ requirement of the opposite sign selection. Only events with zero or one jets are considered. Heavy flavor contributions are effectively removed by the \vec{E}_T requirements. The primary background in this selection region is $W\gamma$ events, where the photon is misidentified as a lepton. CDF uses 4.8 fb^{-1} and analysis technique similar to the described above to set 95% CL upper limits on the production of a SM Higgs boson that are 8.4 times the expected SM cross section[6] (σ_{SM}^H).

4. Tevatron combination and results

Results of both collaborations, in the High and Low mass region, are combined using two different methods: a Bayesian and a modified frequentist technique[9]. Both perform a counting experiment for each bin of the final discriminant, including effects from systematics uncertainties.

This combination procedure is able to correlate systematics uncertainties among different analysis and experiments. Systematics uncertainties are divided in two main categories: rate and shape systematics. The first ones affect the normalization of the different signal and background contributions; these are the most important and are dominated by theoretical uncertainties on signal and background cross sections used to normalize our simulations. Shape systematics affect the shape of the output discriminant; an important example is the jet energy scale calibration.

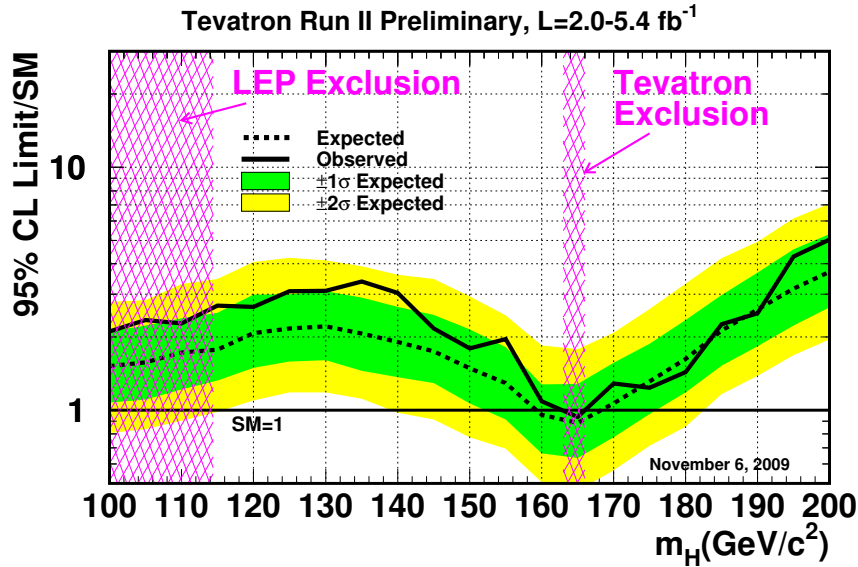


Figure 1: Observed and expected (median, for the background-only hypothesis) 95% CL limits on the ratios to the SM cross section, as functions of the Higgs boson mass for the combined CDF and DØ analyses.

Figure 1 summarize the combination for Higgs masses between 100 and 200 GeV/c^2 . The dotted line represent the median, the green and yellow bands one and two sigma spread of the distribution of the expected limits from a background-only hypothesis; the solid line is the limit that is set looking at data. Each analysis is performed for different Higgs mass hypotheses in 5 GeV steps, then results are connected with a straight line for better readability. The combination **excludes at 95% CL a Standard Model Higgs boson in the $162 < m_H < 166 \text{ GeV}/c^2$ mass range; while the expected exclusion region is $159 < m_H < 169 \text{ GeV}/c^2$.**

5. Conclusions

During latest years CDF and DØ have improved their sensitivity to low cross section processes, are now in the reach of the Higgs production. The combination of the analysis carried out by the two experiments has led to the exclusion of a SM Higgs in the 162 – 166 GeV mass range with expected exclusion range of 159 – 169 GeV. More data is available to be analyzed and more will be collected in the next years allowing an exclusion by each experiment and widening the combined exclusion region.

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