

Measurement of diboson production in early LHC data

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We investigate the discovery potential for the $pp \rightarrow WZ^0 \rightarrow \ell^\pm \nu \ell^+ \ell^-$ ($\ell = e, \mu$) process at CMS, using full detector simulation of the signal and of the various physics and instrumental backgrounds. We propose data driven methods to extract backgrounds and estimate that we can reach 5σ significance of WZ^0 signal at CMS with less than 350 pb^{-1} at 95 % C.L..

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1. Introduction

The study of multiple gauge-boson production at the TeV scale constitutes a unique opportunity to test the Standard Model of Electroweak interactions at the highest possible energies. Production cross sections for all diboson processes, WZ^0 , Z^0Z^0 , WW , $W\gamma$ and $Z^0\gamma$ are relatively large and these processes are expected to be observed already with luminosities of a few hundreds inverse picobarns at the LHC.

A study of WZ^0 production in pp collisions at a center of mass energy of 14 TeV is presented in the following. The production of WZ^0 events allows probing triple gauge-boson couplings and, therefore, non-Abelian gauge symmetry of the Standard Model at energies never before attained. Any deviation of the strength of these couplings from their Standard Model expectations indicates new physics.

We consider the four leptonic final states involving electrons and muons, labeled in the following as $3e$, $2e1\mu$, $2\mu1e$ and 3μ .

We use signal and background Monte Carlo samples that are processed with the full simulation of the CMS detector [1]. The backgrounds are processes with three or more genuine leptons or misidentified jets in the final state. The largest instrumental backgrounds to all four signatures is due to misidentified jets from associated production of jets and heavy bosons and production of top quarks. Among these backgrounds, the main contribution to the WZ^0 signal is from Z^0 +jets processes. The only irreducible backgrounds to the WZ^0 final state are the Z^0Z^0 production with one of the leptons being mis-reconstructed or lost and $Z^0\gamma$ process where a photon converts and produces a genuine electron.

2. Event Selection

A key element for this study is an effective selection of electrons and muons. The selection criteria must have high efficiency for true isolated leptons from heavy boson decays, while at the same time effectively suppressing leptons from heavy quark decays as well as other objects misidentified as leptons.

The identification of electrons is based on matching of a charged track reconstructed in the central tracker with an energy deposition (supercluster) in the electromagnetic calorimeter (ECAL) [2]. Suppression of fake electrons is achieved by applying requirements on the shape of the ECAL energy deposition, on the matching in position and energy between the charged track and the supercluster. Muon candidates are identified by matching a track reconstructed in the muon detectors with a track reconstructed in the central tracker. To further suppress backgrounds originating mainly from genuine muons from heavy quark decays, the significance of the impact parameter in the transverse plane of a muon candidate must also be consistent with a track originating from the primary interaction point. Furthermore, requirements on calorimetric and track isolation are applied for both electron and muon candidates.

Events stemming from the three-lepton final states of WZ^0 production are collected by electron and/or muon triggers. Events are accepted if they contain at least three charged leptons, either electrons or muons, with $p_T > 15\text{ GeV}$ and $|\eta| < 2.5$ (2.4) for electrons (muons). To select WZ^0

candidates we form all possible Z^0 -boson candidates from same-flavour, opposite-charge lepton pairs.

Events are retained if the invariant mass of the Z^0 boson candidate is between 50 and 120 GeV. The event is rejected if a second Z^0 boson candidate is found with an invariant mass within 20 GeV around the nominal Z^0 mass. If two Z^0 boson candidates are found which share one of the leptons, we select the candidate with the invariant mass closest to the nominal Z^0 mass.

After the Z^0 boson candidate is identified, we search for a third lepton in the event to be associated to the W boson decay. This lepton is required to carry a transverse momentum above 20 GeV. If more than one lepton candidate is available, the one with the highest transverse momentum is chosen.

Further suppression of the background can be achieved by applying W identification requirements. We define a W boson candidate transverse mass as

$$M_T(W) = \sqrt{2 \cdot MET \cdot E_\ell (1 - \cos \Delta\phi_{MET,\ell})}, \quad (2.1)$$

where MET is the missing transverse energy, E_ℓ is the energy of the W lepton candidate, and $\Delta\phi_{MET,\ell}$ is the azimuthal separation between the MET direction and the W lepton. We require the transverse mass to be larger than 50 GeV. The event yield for the signal and background events after the full selection is summarized in Table 1.

	all	3e	2e1μ	2μ1e	3μ
WZ^0	34.9±0.5	7.9±0.3	8.0±0.3	8.9±0.3	10.1±0.3
$Z^0 + jets$	3.9±1.0	1.9±0.8	< 0.1	1.8±0.7	0.1±0.1
bbll	2.9±0.3	1.2±0.2	0.1±0.1	1.3±0.2	0.3±0.1
$t\bar{t} + jets$	2.2±0.6	0.6±0.3	0.6±0.3	0.6±0.3	0.3±0.2
$W + jets$	0.4±0.4	0.4±0.4	< 0.1	< 0.1	< 0.1
Z^0Z^0	2.8±0.3	0.8±0.1	0.6±0.1	0.7±0.1	0.7±0.1
$Z^0 + \gamma$	1.4±0.1	0.7±0.1	< 0.1	0.7±0.1	< 0.1
Total non genuine Z^0 bkg	2.6±0.7	1.0 ±0.5	0.6±0.3	0.6±0.3	0.3±0.2
Total genuine Z^0 instrumental bkg	6.8 ±1.0	3.1±0.8	0.1±0.1	3.1±0.5	0.4±0.1
Total genuine Z^0 physics bkg	4.2 ± 0.3	1.5±0.1	0.6±0.1	1.4±0.1	0.7± 0.1

Table 1: Expected number of selected events for an integrated luminosity of 300 pb⁻¹ for the signal and background with 81 GeV < M_Z < 101 GeV obtained using MC truth information. The uncertainties are statistical.

3. Signal Extraction

We separate backgrounds into three categories: physics background from $Z^0\gamma$ and Z^0Z^0 production, processes without a genuine Z^0 boson from $t\bar{t} + jets$ and $W + jets$ production, and processes with a genuine Z^0 boson from $Z^0 + jets$ production.

The physics background from Z^0Z^0 and $Z^0\gamma$ production is estimated from Monte Carlo simulation. We assign a conservative 100% uncertainty on the estimated Z^0Z^0 and $Z^0\gamma$ contribution

due to modeling of kinematics of the decay products and the photon conversion probability. We also use the Monte Carlo simulation to estimate the background from $t\bar{t} + jets$ and $W + jets$ processes, and we assign an additional 100% systematic uncertainty on that contribution. With increasing luminosity, this contribution can be estimated from the side-bands of the Z^0 candidate invariant mass distribution.

The remaining and also the largest instrumental background is from processes with real Z^0 boson and jets where one of the jets is misidentified as a lepton. We estimate this background in data by applying a method, commonly referred to as a “matrix method” (see for example [3]). The idea of the method is to apply a “loose” identification criteria on the lepton associated to the W boson decay, after a Z^0 boson candidate is identified and count the number of observed events, N_{loose} . These events contain events with real isolated leptons N_ℓ and events with misidentified jets N_j :

$$N_{loose} = N_\ell + N_j. \quad (3.1)$$

We then apply a “tight” selection on the third lepton, the number of observed events N_{tight} changes as following

$$N_{tight} = \varepsilon_{tight}N_\ell + p_{fake}N_j, \quad (3.2)$$

where ε_{tight} and p_{fake} are efficiency of “tight” criteria with respect to “loose” requirements for true isolated leptons and misidentified jets, respectively. The “tight” selection uses the final selection criteria applied on the third lepton as used in the analysis, while some of these requirements are loosened for the “loose” sample.

As N_{loose} and N_{tight} are directly observable, to extract the number of $Z + jets$ events in the final sample, one needs to measure ε_{tight} and p_{fake} in control data samples. To estimate ε_{tight} we apply the “tag-and-probe” method [4] using $Z^0 \rightarrow e^+e^-$ or $Z^0 \rightarrow \mu^+\mu^-$ events. The determination of p_{fake} is done on a control-sample of jets misidentified as leptons. The jet composition (heavy vs light flavour, quark vs gluon jets) of the control sample must be similar to the one of the $Z^0 + jets$ events passing the signal selection. Thus, we propose to measure p_{fake} in a $W + jets$ sample which has a similar jet composition as in $Z^0 + jets$. The obtained p_{fake} is further cross-checked with p_{fake} measured in multijet events. We performed both measurements and found the p_{fake} values to agree within uncertainties.

Using the values of ε_{tight} and p_{fake} obtained from the methods described above, we estimated the backgrounds from genuine Z^0 decays by solving Eqs. 3.1 and 3.2 for N_j . The results are given in Table 2. They agree well with the true Monte Carlo yields.

4. Results

The distribution of the Z^0 candidate invariant mass for all four channels combined after applying the final selection is shown in Fig. 1.

We estimated full systematic uncertainties for each of the four individual signature, they amount to 34%, 25%, 35% and 21% for the $3e$, $2e1\mu$, $2\mu1e$ and 3μ final states, respectively. The largest contributions come from the luminosity determination and the requirement on $M_T(W)$. It is comparable to the statistical uncertainty which is roughly 30% for each channel.

	3e	2e1 μ	2 μ 1e	3 μ
$N - ZZ - Z\gamma - W+\text{jets} - t\bar{t}$	11.1 ± 1.3	8.2 ± 0.9	12.1 ± 1.2	10.5 ± 0.8
$N^{\text{genuine } Z}$ (matrix method)	3.2 ± 1.7	0.6 ± 0.8	4.6 ± 2.0	0.6 ± 0.9
N^{WZ^0}	7.9 ± 2.1	7.6 ± 1.2	7.5 ± 2.3	10.0 ± 1.2
WZ^0 from MC	7.9	8.1	9.0	10.1

Table 2: Expected number of events for an integrated luminosity of 300 pb^{-1} for the signal and estimated background for $81 \text{ GeV} < M_Z < 101 \text{ GeV}$ using data-driven methods. Uncertainty is systematic associated with the background subtraction method only.

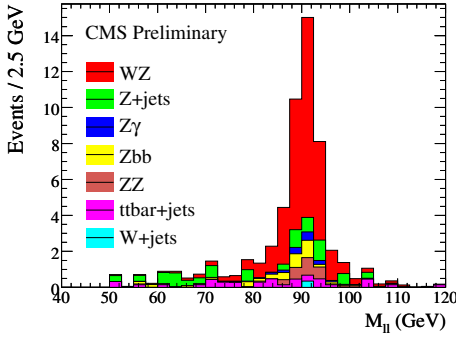


Figure 1: Z^0 candidate invariant mass for all four channels combined, normalized to integrated luminosity of 300 pb^{-1} .

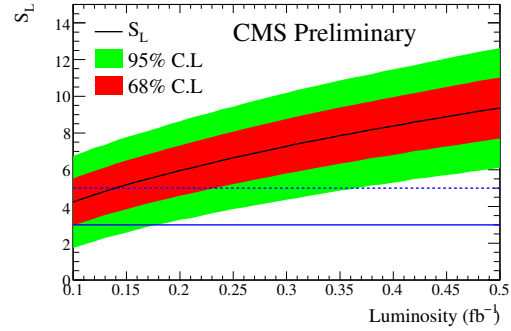


Figure 2: Expected signal significance for WZ^0 production as a function of integrated luminosity. We use a frequentist approach to estimate variation of expected signal and background events. The corresponding 68% and 95% C.L. regions are displayed as red and green bands, respectively.

The expected signal significance S_L of a WZ^0 production as a function of integrated luminosity for all four categories combined, after applying final selection criteria and requirement on the Z^0 invariant mass to be within 10 GeV from the nominal Z^0 boson mass, taking into account full systematic and statistical uncertainties is shown in Fig. 2. We expect to achieve $S_L = 5$ with less than 350 pb^{-1} of integrated luminosity at 95% C.L..

References

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